

# Conformal Renormalization and Energy Functionals in AdS gravity

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HOLOGRAPHYCL

Renormalization of bulk/surface functionals

Minimal Renormalization Scheme: Counterterms

Counterterms and Counterterms in AdS gravity

Beyond Counterterms: Conformal Renormalization

Renormalization of Conformal-2 Functionals

Renormalization of bulk/surface functionals

Renormalization of  $C$ -quasimodes

Counterterms and Counterterms in AdS gravity

Renormalization of  $C$ -quasimodes

Renormalization of  $C$ -quasimodes-2

Renormalization of bulk/surface functionals

Alternative Renormalization Scheme: Kounterterms

Counterterms and Kounterterms in AdS gravity

Beyond Holograms: Conformal Renormalization

Renormalization of Codimension-2 Functionals

Renormalization of bulk/surface functionals

Alternative Renormalization Scheme: Kounterterms

Counterterms and Kounterterms in AdS gravity

Energy Functionals and Renormalization

Renormalization of Chern-Simons? Counterterms

Renormalization of bulk/surface functionals

Alternative Renormalization Scheme: Kounterterms

Counterterms and Kounterterms in AdS gravity

Beyond Einsteinians: Conformal Renormalization

Renormalization of Conformal-2 Functionals

Renormalization of bulk/surface functionals

Alternative Renormalization Scheme: Kounterterms

Counterterms and Kounterterms in AdS gravity

Renormalization of Chern-Simons Gravity

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Renormalization of bulk/surface functionals

Alternative Renormalization Scheme: Kounterterms

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Beyond Kounterterms: Conformal Renormalization

Renormalization of Dimension-2 Functionals



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Renormalization of bulk/surface functionals

Alternative Renormalization Scheme: Kounterterms

Counterterms and Kounterterms in AdS gravity

Beyond Kounterterms: Conformal Renormalization

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Renormalization of Codimension-2 Functionals

Euclidean static black hole metric

$$ds^2 = f^2(r)d\tau^2 + \frac{dr^2}{f^2(r)} + r^2 d\Omega_{D-2}^2, \quad f^2(r) = 1 - \frac{2\omega_D GM}{r^{D-3}} + \frac{r^2}{\ell^2}$$

Einstein-AdS gravity

$$I_{EH} = \frac{1}{16\pi G} \int_M d^D x \sqrt{-g} (R - 2\Lambda), \quad \Lambda = -\frac{(D-1)(D-2)}{2\ell^2}$$

Black hole entropy  $S_{BH} = T I_{EH}^E$

$$T I_{EH}^E = \frac{(D-3)}{(D-2)} M - TS + \lim_{r \rightarrow \infty} \frac{V(S^{D-2})}{8\pi G} \frac{r^{D-1}}{\ell^2}$$

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Gibbs free energy  $G = TI^E$

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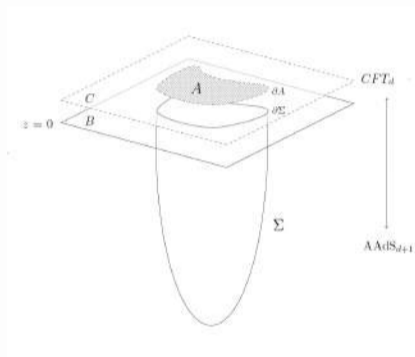
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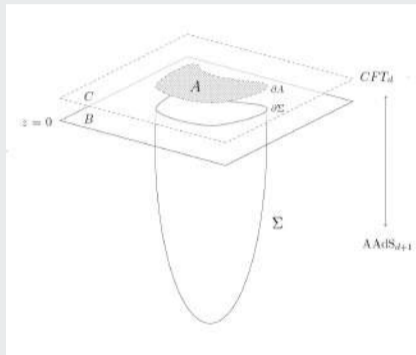
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Holographic Entanglement Entropy from Minimal Surface [Ryu-Takayanagi, 2006]



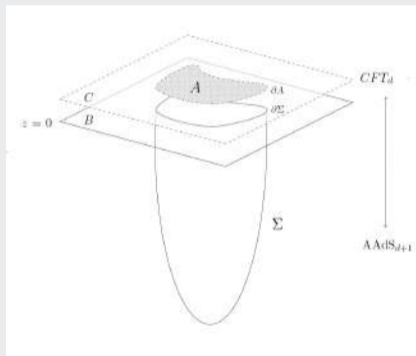
$$S = -\partial_\alpha \lim_{\alpha \rightarrow 1} \mathcal{I}_{\text{grav}}^{(\alpha)}$$

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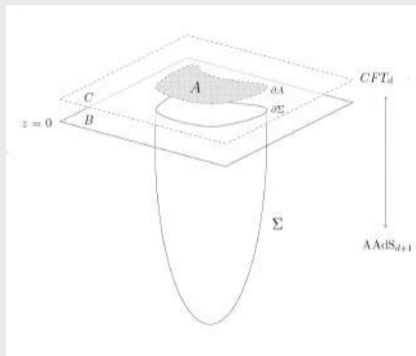
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## HEE from a Cosmic Brane [Lewkowycz-Maldacena, 2013]

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Bulk gravity action evaluated in conical defects

$$I_{\text{EH}}^{(\alpha)} = \frac{1}{16\pi G} \int_{M^{(\alpha)}} d^4x \sqrt{g} R^{(\alpha)} = \frac{1}{16\pi G} \int_M d^4x \sqrt{g} R + \frac{(1-\alpha)}{4G} \mathcal{A}[\Sigma]$$

Area Functional (conformal-2)

$$\mathcal{A}[\Sigma] = \int_{\Sigma} d^2y \sqrt{\gamma}$$

Dual to Einstein gravity

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Holographic Renormalization [Henningson and Skenderis, 1998]

$$I_{\text{ren}} = I_{EH} - \frac{1}{8\pi G} \int_{\partial M} d^d x \sqrt{-h} K + \int_{\partial M} d^d x L_{ct}(h, \mathcal{R}, \nabla \mathcal{R})$$

Counterterms [de Harmsenian-Kraus, 1999], [Kapranov, Johnson, Myers, 1998]

$$8\pi G L_{ct} = \frac{d-1}{l} \sqrt{-h} + \frac{l\sqrt{-h}}{2(d-2)} \mathcal{R} + \frac{l^2 \sqrt{-h}}{2(d-2)^2(d-4)} \left( \mathcal{R}^{ij} \mathcal{R}_{ij} - \frac{d}{4(d-1)} \mathcal{R}^2 \right) \\ + \frac{l^3 \sqrt{-h}}{(d-2)^2(d-4)(d-6)} \left( \frac{8d-2}{3(d-1)} \mathcal{R} \mathcal{R}^{ij} \mathcal{R}_{ij} - \frac{d(d+2)}{16(d-1)^2} \mathcal{R}^3 \right. \\ \left. - 2\mathcal{R}^{ij} \mathcal{R}^{kl} \mathcal{R}_{ijkl} - \frac{d}{3(d-1)} \nabla_i \mathcal{R} \nabla^i \mathcal{R} + \nabla^k \mathcal{R}^{ij} \nabla_k \mathcal{R}_{ij} \right) + \dots$$

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Counterterms [de Boer, Prisco, Skenderis, 1999], [de Boer, Johnson, Myers, 1999]

$$8\pi G L_{ct} = \frac{d-1}{2} \sqrt{-h} + \frac{1}{2(d-2)} \sqrt{-h} \mathcal{R} + \frac{1}{2(d-2)^2(d-4)} \sqrt{-h} \left( \mathcal{R}^{\mu\nu} \mathcal{R}_{\mu\nu} - \frac{d}{4(d-1)} \mathcal{R}^2 \right) \\ + \frac{1}{(d-2)^2(d-4)(d-6)} \sqrt{-h} \left( \frac{d-2}{4(d-1)} \mathcal{R} \mathcal{R}^{\mu\nu} \mathcal{R}_{\mu\nu} - \frac{d(d+2)}{16(d-1)^2} \mathcal{R}^3 \right. \\ \left. - 2\mathcal{R}^{\mu\nu} \mathcal{R}^{\rho\sigma} \mathcal{R}_{\mu\rho} \mathcal{R}_{\nu\sigma} - \frac{d}{4(d-1)} \nabla_\mu \mathcal{R} \nabla^\mu \mathcal{R} + \nabla^\mu \mathcal{R}^\nu \nabla_\mu \mathcal{R}_\nu \right) + \dots$$

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$$\begin{aligned} 8\pi G L_{ct} = & \frac{d-1}{\ell} \sqrt{-h} + \frac{\ell \sqrt{-h}}{2(d-2)} \mathcal{R} + \frac{\ell^3 \sqrt{-h}}{2(d-2)^2(d-4)} \left( \mathcal{R}^{ij} \mathcal{R}_{ij} - \frac{d}{4(d-1)} \mathcal{R}^2 \right) \\ & + \frac{\ell^5 \sqrt{-h}}{(d-2)^3(d-4)(d-6)} \left( \frac{3d-2}{4(d-1)} \mathcal{R} \mathcal{R}^{ij} \mathcal{R}_{ij} - \frac{d(d+2)}{16(d-1)^2} \mathcal{R}^3 \right. \\ & \left. - 2\mathcal{R}^{ij} \mathcal{R}^{kl} \mathcal{R}_{ijkl} - \frac{d}{4(d-1)} \nabla_i \mathcal{R} \nabla^i \mathcal{R} + \nabla^k \mathcal{R}^{ij} \nabla_k \mathcal{R}_{ij} \right) + \dots \end{aligned}$$

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Counterterm Method reproduces BH Thermo

$$G = U - TS$$

Internal Energy

$$U = M + E_0$$

Counter Energy in  $D = 2n + 1$  dimensions

$$E_0 = (-1)^n \frac{(2n-1)!!^2}{(2n)!} \frac{\text{Vol}(S^{2n-1})}{8\pi G} \ell^{2n-2}$$



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## Extrinsic counterterms

$$\tilde{I}_{ren} = I + c_d \int_{\partial M} d^d x B_d(h, K, \mathcal{R})$$

Kounterterms = counterterms of unusual sort (depend on  $K_{ij}$  and  $\mathcal{R}_{ij}^{kl}(h)$ )

$$B_{2n-1} = 2n\sqrt{-h} \int_0^1 dt \delta_{[01 \dots j_{2n-1}]}^{[i_1 \dots i_{2n-1}]} K_{i_1}^{j_1} \left( \frac{1}{2} \mathcal{R}_{i_2 i_3}^{j_2 j_3} - t^2 K_{i_2}^{j_2} K_{i_3}^{j_3} \right) \times \dots$$
$$\dots \times \left( \frac{1}{2} \mathcal{R}_{i_{2n-2} i_{2n-1}}^{j_{2n-2} j_{2n-1}} - t^2 K_{i_{2n-2}}^{j_{2n-2}} K_{i_{2n-1}}^{j_{2n-1}} \right)$$

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$D = 2n$  dimensions [R.O., hep-th/0504233]

$$B_{2n-1} = 2n\sqrt{-h} \int_0^1 dt \delta_{[j_1 \dots j_{2n-1}]^{[i_1 \dots i_{2n-1}]} K_{i_1}^{j_1} \left( \frac{1}{2} \mathcal{R}_{i_2 i_3}^{j_2 j_3} - t^2 K_{i_2}^{j_2} K_{i_3}^{j_3} \right) \times \dots$$
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Kounterterms in  $D = 2n + 1$  [R.O., hep-th/0610230]

$$B_{2n} = 2n\sqrt{-h} \int_0^1 dt \int_0^t ds \delta_{[i_1 \dots i_{2n}] }^{[j_1 \dots j_{2n}]} K_{j_1}^{i_1} \delta_{j_2}^{i_2} \left( \frac{1}{2} \mathcal{R}_{j_3 j_4}^{i_3 i_4} - t^2 K_{j_3}^{i_3} K_{j_4}^{i_4} + \frac{s^2}{\ell^2} \delta_{j_3}^{i_3} \delta_{j_4}^{i_4} \right) \times \dots$$
$$\dots \times \left( \frac{1}{2} \mathcal{R}_{j_{2n-1} j_{2n}}^{i_{2n-1} i_{2n}} - t^2 K_{j_{2n-1}}^{i_{2n-1}} K_{j_{2n}}^{i_{2n}} + \frac{s^2}{\ell^2} \delta_{j_{2n-1}}^{i_{2n-1}} \delta_{j_{2n}}^{i_{2n}} \right) .$$

In  $D = 2n$  dimensions

$$\begin{aligned}\text{tr}(F^n) &= dL_{2n-1}^{CS}(A) \\ F &= dA + A \wedge A\end{aligned}$$

Explicit realization of Chern-Simons forms

$$L_{2n-1}^{CS}(A) = n \int_0^1 dt \text{tr} [A F_t^{n-1}] \quad F_t = tdA + t^2 A^2$$

Global issues (topology)

$$\int_{M_{2n}} (\text{Euler})_{2n} = (4\pi)^n n! \chi(M_{2n}) + \int_{\partial M_{2n}} B_{2n-1}$$

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Global Invariant (Topology)

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## Explicit realization of Chern-Simons forms

$$L_{2n-1}^{CS}(A) = n \int_0^1 dt \mathrm{tr} [A F_t^{n-1}] \quad F_t = tdA + t^2 A^2$$

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Gauge-invariant extension of CS forms  $L_{2n+1}^{TF}(A, \bar{A}) = L_{2n+1}^{CS}(A) - L_{2n+1}^{CS}(\bar{A}) + d\beta_{2n}(A, \bar{A})$

Contact term

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## Black Hole Thermodynamics

$$T I_{bulk}^E = \frac{(D-3)}{(D-2)} M - TS + \lim_{r \rightarrow \infty} \frac{V(S^{D-2})}{8\pi G} \frac{r^{D-1}}{\ell^2}$$

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$$T c_d \int_{\partial M} B_d = \frac{M}{(D-2)} + E_0 - \lim_{r \rightarrow \infty} \frac{V(S^{D-2})}{8\pi G} \frac{r^{D-1}}{\ell^2}$$

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Fefferman-Graham expansion for AAdS Einstein spaces

$$ds^2 = \frac{\ell^2}{z^2} dz^2 + \frac{1}{z^2} g_{ij}(x, z) dx^i dx^j, \quad g_{ij}(x, \rho) = g_{(0)ij}(x) + z^2 g_{(2)ij}(x) + \dots$$

Dirichlet boundary conditions for metrics in AdS spaces

(Lopez-Andreu and Mesterica, 2005)

$$h_{ij} = \frac{g_{(0)ij}}{z^2} + \dots$$

Renormalization = variational problem in  $g_{(0)}$

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Asymptotic form of the extrinsic curvature

$$K_{ij} = \frac{1}{\ell} \frac{g_{(0)ij}}{z^2} + \dots$$

Counterterms of a dual field theory

$$\tilde{I}_{ren} = I_{EH} + c_d \int_{\partial M} d^d x B(f(h), K)$$

as long as the theory is holographic

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AdS gravity action + KTs

$$\tilde{I}_{\text{ren}} = I_{EH} + \frac{\ell^2}{16\pi G} \int_{\partial M} d^3x \sqrt{-h} \delta_{[j_1 j_2 j_3]}^{[i_1 i_2 i_3]} K_{i_1}^{j_1} \left( \frac{1}{2} \mathcal{R}_{i_2 i_3}^{j_2 j_3}(h) - \frac{1}{3} K_{i_2}^{j_2} K_{i_3}^{j_3} \right).$$

Adding zero...

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Counterterms turn into counterterms [O. Miskovic and R.O., 0902.2082]

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Kounterterms turn into counterterms [O.Miskovic and R.O., 0902.2082]

$$L_{ct} = \frac{1}{8\pi G} \frac{\sqrt{-g}}{z^3} \left( \frac{2}{\ell} + \frac{\ell}{2} z^2 \mathcal{R}(g) \right) + \mathcal{O}(z) \\ = \frac{1}{8\pi G} \sqrt{-h} \left( \frac{2}{\ell} + \frac{\ell}{2} \mathcal{R}(h) \right) .$$

EH-AdS gravity +KTs

$$B_{2n-1} = 2n\sqrt{-h} \int_0^1 dt \delta_{[j_1 \dots j_{2n-1}]^{[i_1 \dots i_{2n-1}]} K_{i_1}^{j_1} \left( \frac{1}{2} \mathcal{R}_{i_2 i_3}^{j_2 j_3} - t^2 K_{i_2}^{j_2} K_{i_3}^{j_3} \right) \times \dots$$

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Well...almost. [G.Anastasiou, O.Miskovic, R.O. and I.Papadimitriou, 2003.06425]

$$\tilde{I}_{\text{ren}} = I_{\text{HR}} - \frac{\ell^3}{64\pi G(2n-3)(2n-5)} \int_{\partial M} \sqrt{-h} \mathcal{W}^{ijkl} \mathcal{W}_{ijkl} + \dots$$

A similar result in  $D = 2n + 1$  dimensions.

Last term is zero for most AAdS spaces, but not for gravitational instantons.

Patching up the theory

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Euler-Gauss-Bonnet Theorem in 4D

$$\int_{\partial M} d^3x B_3(K, \mathcal{R}) = \int_M d^4x GB - 32\pi^2 \chi(M)$$

4D Renormalized AdS action [R. Aron et al, gr-qc/9909015]

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GB coupling is also singled out by SUSY [Andrianopoli and D'Auria, arXiv:1405.2010]

Renormalized Einstein-Maxwell-Dirac action (1973)

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Weyl tensor

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$$W_{\mu\nu}^{\alpha\beta} = R_{\mu\nu}^{\alpha\beta} - 4S_{[\mu}^{\alpha} \delta_{\nu]}^{\beta]}, \quad \text{Schouten } S_{\mu}^{\alpha} = \frac{1}{D-2} (R_{\mu}^{\alpha} - \frac{1}{2(D-1)} \delta_{\mu}^{\alpha} R)$$

GB coupling is also singled out by SUSY [Andrianopoli and D'Auria, arXiv:1405.2010]

Renormalized AdS action = MacDowell-Mansouri action (1977)

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Renormalized action for Einstein spaces

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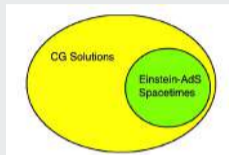
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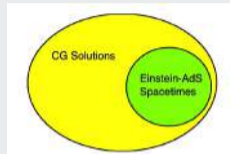
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## Embedding Einstein theory in Conformal Gravity



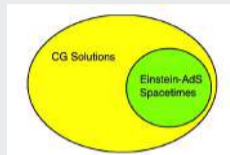
- Why?: Conformal Gravity is finite for AAdS conditions. [Grumiller et al., 2013]
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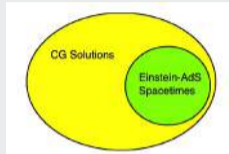
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Einstein gravity from CG with Neumann bc's [Maldacena, 2011]

$$I_{CG} = \alpha_{CG} \int_M d^4x \sqrt{-g} W_{\mu\nu\alpha\beta} W^{\mu\nu\alpha\beta}$$

Foliation-Graham expansion for AdS spaces in CG

$$ds^2 = \frac{\ell^2}{z^2} dz^2 + \frac{1}{z^2} g_{ij}(x, z) dx^i dx^j, \quad g_{ij}(x, z) = g_{(0)ij}(x) + z^2 g_{(2)ij}(x) + \dots \\ + z g_{(1)ij}(x) + \dots$$

Weyl Conformal Gravity

$$B_{\mu\nu} = \nabla^\lambda C_{\mu\nu\lambda} + S^{\lambda\sigma} W_{\lambda\mu\nu\sigma} = 0, \quad C_{\mu\lambda}^\sigma = \nabla_\nu S_{\lambda}^\sigma - \nabla_\lambda S_{\nu}^\sigma$$

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[Imbimbo, Schwimmer, Theisen and Yankielowicz, hep-th/9910267]

CG action for Einstein spaces = Renormalized Einstein-AdS action

$$I_{CG}[E] = I_{HR}$$

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EH Action+Euler term

$$\tilde{I}_{\text{ren}} = \frac{1}{16\pi G} \int_M d^6x \sqrt{-g} \left( R + \frac{20}{\ell^2} - \frac{\ell^4}{72} (\text{Euler})_6 \right),$$

In terms of fully-antisymmetric objects

$$\tilde{I}_{\text{ren}} = \frac{1}{16\pi G \times 192} \int_M d^6x \sqrt{-g} \delta_{[\mu_1 \dots \mu_3]}^{[\nu_1 \dots \nu_3]} \left[ R_{\nu_1 \nu_2}^{\mu_1 \mu_2} \delta_{[\mu_3 \mu_4]}^{[\nu_3 \nu_4]} \delta_{[\mu_5 \mu_6]}^{[\nu_5 \nu_6]} \right. \\ \left. + \frac{2}{3/2} \delta_{[\mu_1 \mu_2]}^{[\nu_1 \nu_2]} \delta_{[\mu_3 \mu_4]}^{[\nu_3 \nu_4]} \delta_{[\mu_5 \mu_6]}^{[\nu_5 \nu_6]} - \frac{\ell^4}{3} R_{\nu_1 \nu_2}^{\mu_1 \mu_2} R_{\nu_3 \nu_4}^{\mu_3 \mu_4} R_{\nu_5 \nu_6}^{\mu_5 \mu_6} \right],$$

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Polynomial of  $W_{(E)}$

$$\begin{aligned} \tilde{I}_{\text{ren}} = & \frac{\ell^4}{16\pi G \times 4!} \int_M d^6 x \sqrt{-g} \left[ \frac{1}{2\ell^2} \delta_{[\mu_1 \dots \mu_4]}^{[\nu_1 \dots \nu_4]} W_{(E)\nu_1\nu_2}^{\mu_1\mu_2} W_{(E)\nu_3\nu_4}^{\mu_3\mu_4} \right. \\ & \left. - \frac{1}{4!} \delta_{[\mu_1 \dots \mu_6]}^{[\nu_1 \dots \nu_6]} W_{(E)\nu_1\nu_2}^{\mu_1\mu_2} W_{(E)\nu_3\nu_4}^{\mu_3\mu_4} W_{(E)\nu_5\nu_6}^{\mu_5\mu_6} \right], \end{aligned}$$



There are three Conformal Invariants in 6D

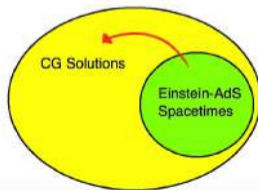
$$\begin{aligned}I_1 &= W_{\alpha\beta\mu\nu} W^{\alpha\sigma\lambda\nu} W_{\sigma}{}^{\beta\mu}{}_{\lambda}, \\I_2 &= W_{\mu\nu\alpha\beta} W^{\alpha\beta\sigma\lambda} W_{\sigma\lambda}{}^{\mu\nu}, \\I_3 &= W_{\mu\rho\sigma\lambda} \left( \delta_{\nu}^{\mu} \square + 4R_{\nu}^{\mu} - \frac{6}{5}R\delta_{\nu}^{\mu} \right) W^{\nu\rho\sigma\lambda} + \nabla_{\mu} J^{\mu},\end{aligned}$$

with

$$\begin{aligned}J_{\mu} &= 4R_{\mu}{}^{\lambda\rho\sigma} \nabla^{\nu} R_{\nu\lambda\rho\sigma} + 3R^{\nu\lambda\rho\sigma} \nabla_{\mu} R_{\nu\lambda\rho\sigma} - 5R^{\nu\lambda} \nabla_{\mu} R_{\nu\lambda} \\&\quad + \frac{1}{2}R \nabla_{\mu} R - R_{\mu}^{\nu} \nabla_{\nu} R + 2R^{\nu\lambda} \nabla_{\nu} R_{\lambda\mu}.\end{aligned}$$

Einstein	→	Conformal	CI's
$\delta_{[\mu_1 \dots \mu_6]}^{[\nu_1 \dots \nu_6]} \mathbf{W}_{(E)\nu_1\nu_2}^{\mu_1\mu_2} \mathbf{W}_{(E)\nu_3\nu_4}^{\mu_3\mu_4} \mathbf{W}_{(E)\nu_5\nu_6}^{\mu_5\mu_6}$	→	$\delta_{[\mu_1 \dots \mu_6]}^{[\nu_1 \dots \nu_6]} \mathbf{W}_{\nu_1\nu_2}^{\mu_1\mu_2} \mathbf{W}_{\nu_3\nu_4}^{\mu_3\mu_4} \mathbf{W}_{\nu_5\nu_6}^{\mu_5\mu_6}$	32(2I <sub>1</sub> + I <sub>2</sub> )
$-\frac{1}{\ell^2} \delta_{[\mu_1 \dots \mu_4]}^{[\nu_1 \dots \nu_4]} \mathbf{W}_{(E)\nu_1\nu_2}^{\mu_1\mu_2} \mathbf{W}_{(E)\nu_3\nu_4}^{\mu_3\mu_4}$	→	$\delta_{[\mu_1 \dots \mu_5]}^{[\nu_1 \dots \nu_5]} \mathbf{W}_{\nu_1\nu_2}^{\mu_1\mu_2} \mathbf{W}_{\nu_3\nu_4}^{\mu_3\mu_4} \mathbf{S}_{\nu_5}^{\mu_5} + 16\mathbf{C}^{\mu\nu\lambda} \mathbf{C}_{\mu\nu\lambda} + \nabla^\mu \mathbf{J}_\mu$	4I <sub>1</sub> - I <sub>2</sub> - I <sub>3</sub>

$$\mathbf{J}_\mu = 16\mathbf{W}_\mu^{\kappa\lambda\nu} \mathbf{C}_{\kappa\lambda\nu} - 2\mathbf{W}_{\nu\sigma}^{\kappa\lambda} \nabla_\mu \mathbf{W}_{\kappa\lambda}^{\nu\sigma}$$



6D CG with an Einstein sector [Lu, Pang and Pope, 2013]

$$I_{CG} = \alpha_{CG} \int_M d^6 x \sqrt{-\hat{g}} \left( \frac{1}{4!} \delta_{[\mu_1 \dots \mu_6]}^{[\nu_1 \dots \nu_6]} W_{\nu_1 \nu_2}^{\mu_1 \mu_2} W_{\nu_3 \nu_4}^{\mu_3 \mu_4} W_{\nu_5 \nu_6}^{\mu_5 \mu_6} + \frac{1}{2} \delta_{[\mu_1 \dots \mu_5]}^{[\nu_1 \dots \nu_5]} W_{\nu_1 \nu_2}^{\mu_1 \mu_2} W_{\nu_3 \nu_4}^{\mu_3 \mu_4} S_{\nu_5}^{\mu_5} \right. \\ \left. + 8C^{\mu\nu\lambda} C_{\mu\nu\lambda} \right) + \alpha_{CG\partial M} d^5 x \sqrt{-h} n^\mu \left( 8W_{\mu}^{\kappa\lambda\nu} C_{\kappa\lambda\nu} - W_{\nu\sigma}^{\kappa\lambda} \nabla_{\mu} W_{\kappa\lambda}^{\nu\sigma} \right).$$

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- Variation of  $I_{CG}$  gives EOM in terms of Weyl, Cotton and Schouten tensors.
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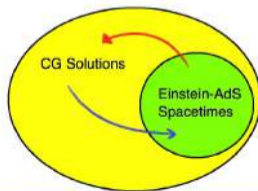
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LPP CG action decomposed into Einstein and non-Einstein parts:

$$I_{CG} = -4! \alpha_{CG} \int_M d^6 x \sqrt{-g} [P_6 (W_{(E)}) + Q (W_{(E)}, H)] \\ - \alpha_{CG} \int_{\partial M} d^5 x \sqrt{-h} n^\mu \left( W_{(E)\nu\sigma}^{\kappa\lambda} \nabla_\mu W_{(E)\kappa\lambda}^{\nu\sigma} \right).$$





# Back to Einstein gravity (with an extra term)



Einstein condition, and  $\alpha_E = -\frac{\ell^4}{384\pi G}$ :

$$I_{CG}[E] = \frac{1}{16\pi G} \int_M d^6x \sqrt{-g} \left( R + \frac{20}{\ell^2} - \frac{\ell^4}{72} (Euler)_6 \right) + \frac{\ell^4}{384\pi G} \int_{\partial M} d^5x \sqrt{-h} n^\mu \left( W_{(E)\nu\sigma}^{\kappa\lambda} \nabla_\mu W_{(E)\kappa\lambda}^{\nu\sigma} \right),$$

Performing asymptotic expansions

$$\Delta I = \frac{\ell^4}{192\pi G} \int_{\partial M} d^5x \sqrt{-h} W^{IJMN}(h) W_{IJMN}(h) + \dots$$

CG action for Einstein spaces = Renormalized Einstein-AdS action

$$I_{CG}[E] = I_{\text{ren}}$$

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CG action for Einstein spaces = Renormalized Einstein-AdS action

$$I_{CG}[E] = I_{HR}$$

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CG action for Einstein spaces = Renormalized Einstein-AdS action

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## Riemann squared term

$$\int_{M^{(\alpha)}} d^4x \sqrt{g} \left( \text{Rie}^{(\alpha)} \right)^2 = \int_M d^4x \sqrt{g} \text{Rie}^2 + 8\pi(1-\alpha) \int_{\Sigma} d^2y \sqrt{\gamma} \left( R_{ABAB} - \mathcal{K}_{mn}^{(A)} \mathcal{K}_{(A)}^{mn} \right) + \dots$$

## Ricci squared term

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## SCALAR SQUARED TERM

$$\int_{M^{(\alpha)}} d^4x \sqrt{g} \left( R^{(\alpha)} \right)^2 = \int_M d^4x \sqrt{g} R^2 + 8\pi(1-\alpha) \int_{\Sigma} \sqrt{\gamma} R + \dots$$

[Fursaev-Patrushev-Solodukhin, 2013]

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## 4D Conformal Gravity

$$I_{CG} = \frac{\ell^2}{64\pi G} \int_M d^4x \sqrt{-g} W_{\mu\nu\alpha\beta} W^{\mu\nu\alpha\beta} - \frac{\pi\ell^2}{2G} \chi[M]$$

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## In a manifold with a conical singularity

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$$\int_{M^{(0)}} d^4x \sqrt{g} |W^{(0)}|^2 = \int_M d^4x \sqrt{g} |W|^2 + 8\pi(1-\alpha) \int_{\Sigma} d^2v \sqrt{\gamma} K_{\Sigma} + \mathcal{O}((1-\alpha)^2)$$

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## 4D Conformal Gravity

$$I_{CG} = \frac{\ell^2}{64\pi G} \int_M d^4x \sqrt{g} \left( Rie^2 - 2Ric^2 + \frac{1}{3}R^2 \right) - \frac{\pi\ell^2}{2G} \chi[M]$$

## In a manifold with a conical singularity

$$\int_{M^{(\alpha)}} d^4x \sqrt{g} |W^{(\alpha)}|^2 = \int_M d^4x \sqrt{g} |W|^2 + 8\pi(1-\alpha) \int_{\Sigma} d^2y \sqrt{\gamma} K_{\Sigma} + \mathcal{O}\left((1-\alpha)^2\right)$$

Conformal Invariant in codimension-2 (Graham-Witten anomaly)

$$K_{\Sigma} = W_{mn}^{mn} - P_{mn}^{(A)} P_{(A)}^{mn},$$

Traceless part of the extrinsic curvature

$$P_{mn}^{(A)} = \mathcal{K}_{mn}^{(A)} - \frac{1}{2} \mathcal{K}^{(A)} \gamma_{mn}$$

Conformal Gravity Action

$$I_{CG}^{(\alpha)} = I_{CG} + \frac{(1-\alpha)}{4G} L_{\Sigma} + \mathcal{O}((1-\alpha)^2)$$

Conformal Invariance in Codimension-2 is inherited from the Bulk

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Traceless part of the extrinsic curvature

$$P_{mn}^{(A)} = \mathcal{K}_{mn}^{(A)} - \frac{1}{2} \mathcal{K}^{(A)} \gamma_{mn}$$

Conformal Gravity Action

$$I_{CS}^{(d)} = I_{CS} + \frac{(1-\alpha)}{4G} L_{\Sigma} + \mathcal{O}((1-\alpha)^2)$$

Conformal Invariance in Codimension-2 is inherited from the Bulk

Conformal Invariant in codimension-2 (Graham-Witten anomaly)

$$K_{\Sigma} = W_{mn}^{mn} - P_{mn}^{(A)} P_{(A)}^{mn} ,$$

Traceless part of the extrinsic curvature

$$P_{mn}^{(A)} = \mathcal{K}_{mn}^{(A)} - \frac{1}{2} \mathcal{K}^{(A)} \gamma_{mn}$$

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$$I_{\text{CG}}^{(\alpha)} = I_{\text{CG}} + \frac{(1-\alpha)}{4G} L_{\Sigma} + \mathcal{O}\left((1-\alpha)^2\right)$$

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Conformal Invariance in Codimension-2 is inherited from the Bulk!

Minimal surface/Einstein ambient space

$$\mathcal{A}_{\text{ren}} = L_{\Sigma_{\text{min}}}[E]$$

Renormalized Area

$$\mathcal{A}_{\text{ren}}[\Sigma] = \frac{\ell^2}{2} \int_{\Sigma} d^2y \sqrt{\gamma} \left[ W_{(E)mn}^{\text{min}} - P_{mn}^{(A)} P_{(A)}^{\text{min}} \right] - 2\pi\ell^2 \chi[\Sigma]$$

Ren. Area [Alexakis-Noronha, 2010] / Ren. HEE [Alexakis-Prays-NO, 2010]

$$\mathcal{A}_{\text{ren}}[\Sigma] = \frac{\ell^2}{4} \int_{\Sigma} d^2y \sqrt{\gamma} \delta_{\text{min}}^{spq} \left( R_{pq}^{\text{min}} + \frac{1}{\ell^2} \delta_{pq}^{\text{min}} \right) - 2\pi\ell^2 \chi[\Sigma]$$

Minimal surface/Einstein ambient space

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see Area [Lopez-Vivanco, 2010] / Ren. HEE [Anastasiou-Argyres-00, 2010]

$$\mathcal{A}_{\text{ren}}[\Sigma] = \frac{\ell^2}{4} \int_{\Sigma} d^2y \sqrt{\gamma} \delta_{\text{min}}^{spq} \left( R_{pq}^{\text{min}} + \frac{1}{\ell^2} \delta_{pq}^{\text{min}} \right) - 2\pi \ell^2 \chi[\Sigma]$$

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Gen. Area [Alexakis-Noronha, 2010] / Ren. HEE [Anastasiou-Argyres-NO, 2010]

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See also [Liu and Y. Wang, 2010] / Pen. HEE [Liu and Y. Wang, 2010]

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Ren. Area [Alexakis-Mazzeo, 2010] / Ren. HEE [Anastasiou-Araya-R0, 2018]

$$\mathcal{A}_{\text{ren}}[\Sigma] = \frac{\ell^2}{4} \int_{\Sigma} d^2y \sqrt{\gamma} \delta_{mn}^{\rho q} \left( \mathcal{R}_{pq}^{mn} + \frac{1}{\ell^2} \delta_{pq}^{mn} \right) - 2\pi\ell^2 \chi[\Sigma]$$



Minimal surface/Einstein ambient space

$$\mathcal{A}_{\text{ren}} = L_{\Sigma_{\text{min}}}[E]$$

Renormalized Area

$$\mathcal{A}_{\text{ren}}[\Sigma] = \frac{\ell^2}{2} \int_{\Sigma} d^2y \sqrt{\gamma} \left[ W_{(E)mn}^{mn} - P_{mn}^{(A)} P_{(A)}^{mn} \right] - 2\pi\ell^2 \chi[\Sigma]$$

Ren. Area [Alexakis-Mazzeo, 2010] / Ren. HEE [Anastasiou-Araya-R0, 2018]

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Functional defined on compact and orientable 2D surfaces immersed in  $\mathbb{R}^3$



In terms of the (spatial) mean curvature

$$W[\Sigma] = \int_{\Sigma} H^2 \sqrt{|g|}^2$$

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$$\mathcal{W}[\Sigma] = \int_{\Sigma} d^2y \sqrt{\gamma} H^2$$

Functional defined on compact and orientable 2D surfaces immersed in  $\mathbb{R}^3$



In terms of the (spatial) mean curvature

$$\mathcal{W}[\Sigma] = \int_{\Sigma} d^2y \sqrt{\gamma} H^2$$

## Renormalized Area

$$\mathcal{A}_{\text{ren}}[\Sigma] = \frac{\ell^2}{2} \int_{\Sigma} d^2y \sqrt{\gamma} \left( W_{(E)mn}^{mn} - P_{mn}^{(A)} P_{(A)}^{mn} \right) - 2\pi\ell^2 \chi[\Sigma]$$

For pure/global  $\text{AdS}_4$  as ambient space, constant-time slice

$$W = 0 \qquad \mathcal{K}^{(t)} = 0$$

## Case: Conformal relations

$$\mathcal{A}_{\text{ren}}[\Sigma] = -\frac{\ell^2}{2} \int_{\Sigma} d^2y \sqrt{\gamma} \left( R_{ij}^ij - \mathcal{R} + 2H^2 \right) - 2\pi\ell^2 \chi[\Sigma]$$

## Renormalized Area

$$\mathcal{A}_{\text{ren}}[\Sigma] = \frac{\ell^2}{2} \int_{\Sigma} d^2y \sqrt{\gamma} \left( W_{(E)mn}^{mn} - P_{mn}^{(A)} P_{(A)}^{mn} \right) - 2\pi\ell^2 \chi[\Sigma]$$

For pure/global  $\text{AdS}_4$  as ambient space, constant-type slice

$$W = 0$$

$$K^{(A)} = 0$$

Using Einstein relations

$$\mathcal{A}_{\text{ren}}[\Sigma] = -\frac{\ell^2}{2} \int_{\Sigma} d^2y \sqrt{\gamma} \left( R_{(E)}^E - \mathcal{R} + 2H^2 \right) - 2\pi\ell^2 \chi[\Sigma]$$

## Renormalized Area

$$\mathcal{A}_{\text{ren}}[\Sigma] = \frac{\ell^2}{2} \int_{\Sigma} d^2y \sqrt{\gamma} \left( W_{(E)mn}^{mn} - P_{mn}^{(A)} P_{(A)}^{mn} \right) - 2\pi\ell^2 \chi[\Sigma]$$

For pure/global  $\text{AdS}_4$  as ambient space, constant-time slice

$$W = 0$$

$$\mathcal{K}^{(t)} = 0$$

Using Einstein relations

$$\mathcal{A}_{\text{ren}}[\Sigma] = -\frac{\ell^2}{2} \int_{\Sigma} d^2y \sqrt{\gamma} \left( R_{(3)}^y{}_y - \mathcal{R} + 2H^2 \right) - 2\pi\ell^2 \chi[\Sigma]$$



Renormalized Area

$$\mathcal{A}_{\text{ren}}[\Sigma] = \frac{\ell^2}{2} \int_{\Sigma} d^2y \sqrt{\gamma} \left( W_{(E)mn}^{mn} - P_{mn}^{(A)} P_{(A)}^{mn} \right) - 2\pi\ell^2 \chi[\Sigma]$$

For pure/global AdS<sub>4</sub> as ambient space, constant-time slice

$$W = 0$$

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For pure/global AdS<sub>4</sub> as ambient space, constant-time slice

$$W = 0$$

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## Gauss-Codazzi relations

$$\mathcal{A}_{\text{ren}}[\Sigma] = -\frac{\ell^2}{2} \int_{\Sigma} d^2y \sqrt{\gamma} \left( R_{ij}^{ij} - \mathcal{R} + 2H^2 \right) - 2\pi\ell^2 \chi[\Sigma]$$

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## Gauss-Codazzi relations

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In the conformal frame  $\hat{g}_{\mu\nu}$

$$\mathcal{A}_{\text{ren}}[\Sigma] = \frac{\ell^2}{2} \int_{\Sigma} d^2y \sqrt{\hat{\gamma}} (\hat{\mathcal{R}} - 2\hat{H}^2) - 2\pi\ell^2 \chi[\Sigma]$$

For a compact surface

$$\int_{\Sigma_{\text{comp}}} d^2y \sqrt{\hat{\gamma}} \hat{\mathcal{R}} = 4\pi \chi[\Sigma_{\text{comp}}]$$

Willmore Energy [Adams-Lazarus, Borsoi, HÜ, Riviere-Bellocquer, 2020]

$$\mathcal{A}_{\text{ren}}[\Sigma_{\text{comp}}] = -\ell^2 \mathcal{W}[\Sigma_{\text{comp}}]$$

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Willmore Energy (AdS/CFT, Maldacena, Witten, Nester, Deser, '92)

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Willmore Energy (Adams, 1974; Borner, 1977; Alvarez-Gaume, 1984)

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$$\mathcal{A}_{\text{ren}}[\Sigma_{\text{comp}}] = -\ell^2 \mathcal{W}[\Sigma_{\text{comp}}]$$

Arbitrary  $\Sigma$ , Einstein ambient space

$$L_{\Sigma}[E] = \frac{\ell^2}{4} I_H[\Sigma] - 2\pi\ell^2 \chi[\Sigma]$$

Reduced Hawking Mass  $I_H$  [Fischetti and Gleason, 2015]

$$I_H[\Sigma] = 2 \int_{\Sigma} d^2y \sqrt{\gamma} \left[ \mathcal{R} + \frac{2}{\ell^2} - \frac{1}{2} (\mathcal{K}^{(A)})^2 \right]$$

Conformal Renormalized Area Functional

$$L_{\Sigma}[E] = \mathcal{A}_{\text{ren}}[\Sigma] - \frac{\ell^2}{4} \int_{\Sigma} d^2y \sqrt{\gamma} (\mathcal{K}^{(A)})^2$$

Arbitrary  $\Sigma$ , Einstein ambient space

$$L_{\Sigma} [E] = \frac{\ell^2}{4} I_H [\Sigma] - 2\pi\ell^2 \chi [\Sigma]$$

Reduced Hawking Mass  $L_{\Sigma} [E]$  (Fischler and Strominger, 2004)

$$I_H [\Sigma] = 2 \int_{\Sigma} d^p y \sqrt{\gamma} \left[ R + \frac{2}{\ell^2} - \frac{1}{2} (\mathcal{K}^{(A)})^2 \right]$$

Reduced Renormalized Area Functional

$$L_{\Sigma} [E] = \mathcal{A}_{ren} [\Sigma] - \frac{\ell^2}{4} \int_{\Sigma} d^p y \sqrt{\gamma} (\mathcal{K}^{(A)})^2$$

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Generalizes Renormalized Area Functional

$$L_{\Sigma} [E] = \mathcal{A}_{\text{ren}} [\Sigma] - \frac{\ell^2}{4} \int_{\Sigma} d^2y \sqrt{\gamma} \left( \mathcal{K}^{(A)} \right)^2$$

Arbitrary  $\Sigma$ , Einstein ambient space

$$L_{\Sigma} [E] = \frac{\ell^2}{4} I_H [\Sigma] - 2\pi\ell^2 \chi [\Sigma]$$

Reduced Hawking Mass  $I_H$  [Fischetti and Wiseman, 2016]

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Generalizes Renormalized Area Functional

$$L_{\Sigma} [E] = \mathcal{A}_{\text{ren}} [\Sigma] - \frac{\ell^2}{4} \int_{\Sigma} d^2y \sqrt{\gamma} \left( \mathcal{K}^{(A)} \right)^2$$

Conformal Invariance in the Bulk  $\implies$  Conformal Invariance in Codimension-2

Renormalization in the Bulk  $\implies$  Renormalization in Codimension 2

Conformal Invariance  $\implies$  Renormalization (1)

Renormalized Volume  $\implies$  Renormalized Area (in conically singular manifolds)

Conformal Energy in higher dimensions  $\implies$   $\mathcal{E}$  (in  $d > 4$  spacetimes)



Conformal Invariance in the Bulk  $\implies$  Conformal Invariance in Codimension-2

Renormalization in the Bulk  $\implies$  Renormalization in Codimension-2

Conformal Anomalies in the Bulk

Renormalized Volume  $\implies$  Renormalized Area for convex regular manifolds

Conformal Energy in higher dimensions  $\implies$  Conformal Anomaly

Conformal Invariance in the Bulk  $\implies$  Conformal Invariance in Codimension-2

Renormalization in the Bulk  $\implies$  Renormalization in Codimension-2

Conformal Invariance  $\implies$  Renormalization (1)

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Conformal Invariance in the Bulk  $\implies$  Conformal Invariance in Codimension-2

Renormalization in the Bulk  $\implies$  Renormalization in Codimension-2

Renormalized Volume  $\implies$  Renormalized Area in gravity coupled matter

Renormalized Energy  $\implies$  Renormalized Energy in gravity coupled matter

Conformal Invariance in the Bulk  $\implies$  Conformal Invariance in Codimension-2

Renormalization in the Bulk  $\implies$  Renormalization in Codimension-2

Conformal Invariance  $\implies$  Renormalization (???)

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Conformal Invariance  $\implies$  Renormalization (???)

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Conformal Gravity in higher even dimensions  $D \geq 8$ ? (with N.Boulanger)



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# Acknowledgements

