

# The Higgs Portal to Cosmology

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**Central idea:**

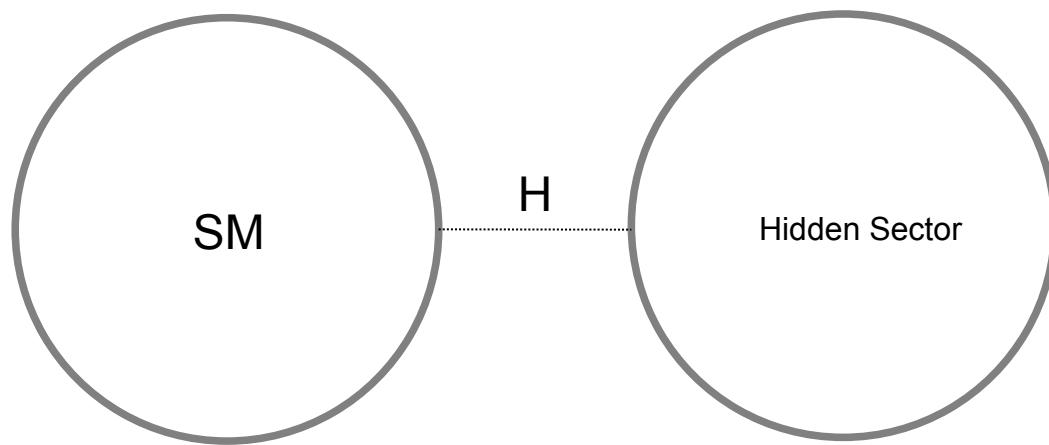
$H^+H^-$  \* ( STUFF )

**Keywords:**

*Higgs, inflation, vacuum stability, dark matter*

( review: [2104.03342 \[hep-ph\]](#) )

# The Higgs and the dark sector



Lowest order operators ("Higgs Portal") :

$$\bar{H}H S^2 + \dots \quad (\text{scalar})$$

$$\bar{H}H V_\mu V^\mu + \dots \quad (\text{vector})$$

$$\bar{H}H \bar{\chi}\chi / \Lambda + \dots \quad (\text{fermion})$$

"Portal" due to Patt, Wilczek'06 (earlier : Silveira, Zee'85 ; Shabinger, Wells'05 ; ...)

Special role of the Higgs :

$|H|^2$  = the only gauge and Lorentz-inv. dim-2 operator

$$L = a |H|^2 S^2 + b |H|^2 S$$

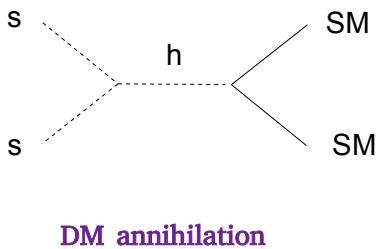
(  $S$  = "hidden" scalar )

$b=0$  ( $S$  has hidden charge):

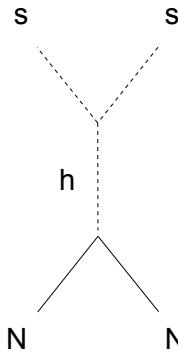
$$L = a |H|^2 S^2$$

" $S$ " is stable and couples weakly to SM  **DARK MATTER (?)**

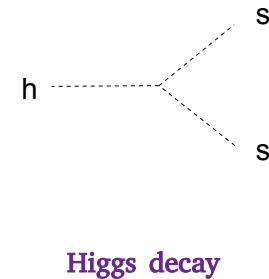
## Dark matter:



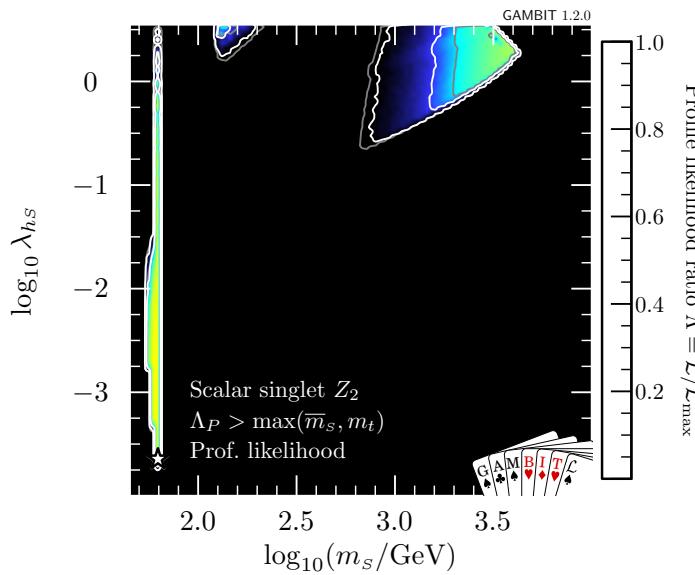
DM annihilation



DM direct detection



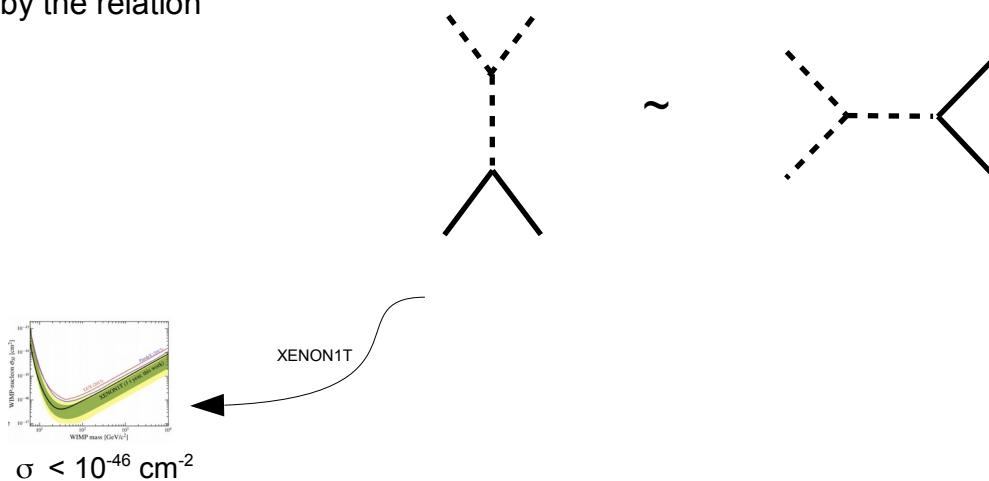
Higgs decay





## *preference for multi-TeV DM masses*

Driven by the relation



**Model dependent. Does not apply if:**

- amplitude vanishes at  $q = 0$
- $\exists$  light unstable dark states
- broad resonant annihilation

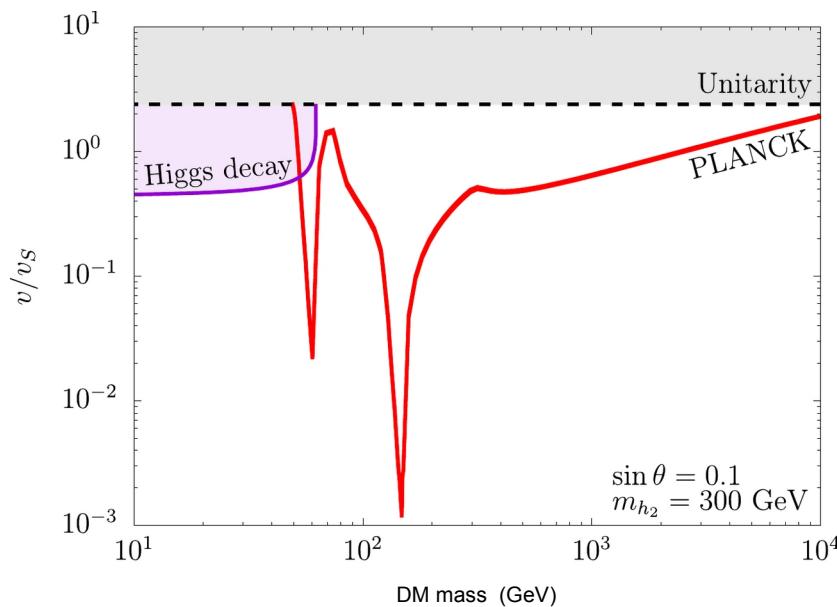
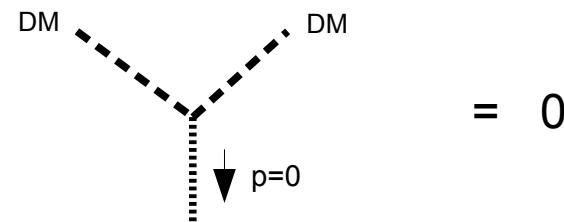
Example: **SM + a complex scalar S** (*pseudo-Goldstone DM*)

$$V = V_0 + V_{\text{soft}},$$

$$V_0 = -\frac{\mu_H^2}{2} |H|^2 - \frac{\mu_S^2}{2} |S|^2 + \frac{\lambda_H}{2} |H|^4 + \lambda_{HS} |H|^2 |S|^2 + \frac{\lambda_S}{2} |S|^4, \quad V_{\text{soft}} = -\frac{\mu'_S^2}{4} S^2 + \text{h.c.}$$

$\text{Im } S = \text{dark matter}$

**Goldstone feature:**



Direct detection = loop-suppressed

Annihilation = unsuppressed



**excellent WIMP**

(from 60 GeV to 10 TeV)

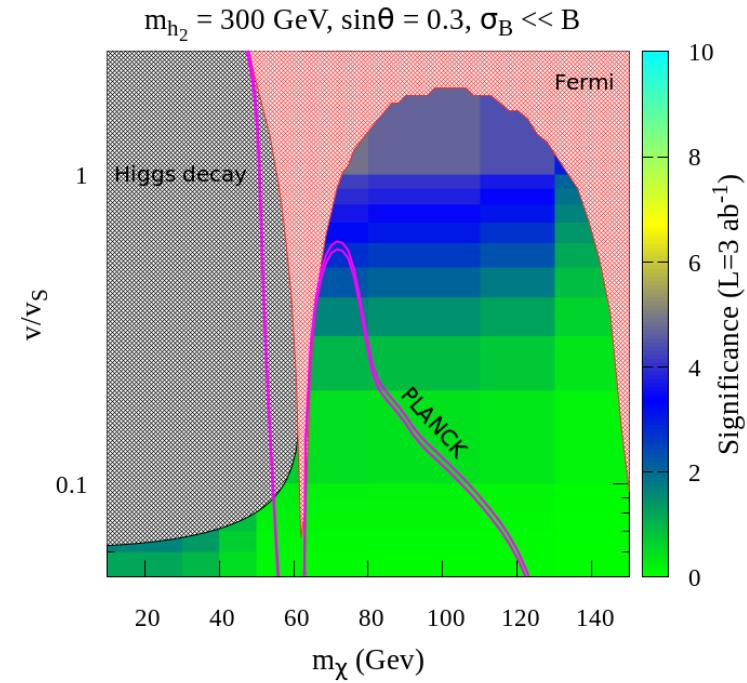
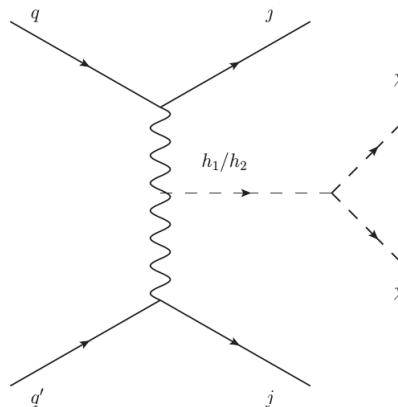
LHC:

DM = missing  $E_T$

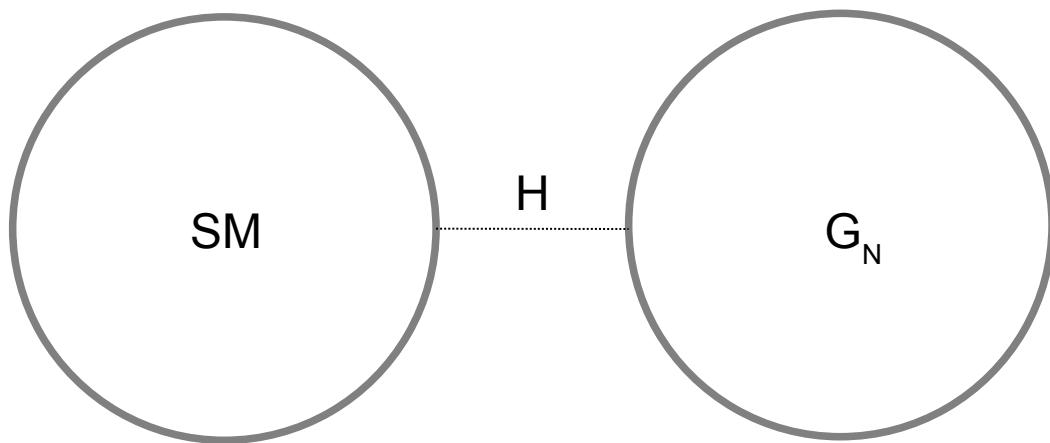
**Higgs portal:** Higgs–mediated production (monojets, VBF, etc.)

**Problem:** limited kinematic reach ( $m_{DM} < m_h/2$ )

Huitu, Koivunen, OL, Mondal '18

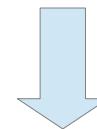


## The Higgs and vector dark matter



$V \sim \bar{H}H \bar{S}S$        $\Rightarrow$       H-S mixing       $\Rightarrow$       **h couples to  $G_N$**

Lie groups possess discrete symmetries



gauge fields as dark matter

E.g.  $U(1) : A_\mu \rightarrow -A_\mu$

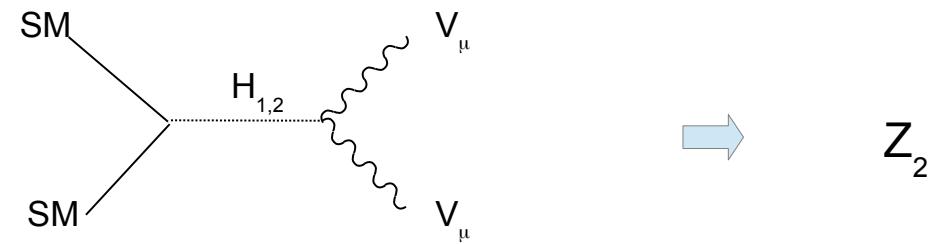
## Higgs mechanism in the hidden sector :

$$\mathcal{L} = -1/4 F_{\mu\nu} F^{\mu\nu} + D_\mu S^* D^\mu S - V(S) + \lambda/4 \bar{H} H S^* S$$

$S \longrightarrow VEV$



SM couplings:



gauge invariance (+ minimal field content)



$Z_2$



gauge fields are natural DM candidates

## SU(N):

Break to nothing in minimal way  $\rightarrow$  N-1 dark-Higgs N-plets

$$\begin{array}{c} \left[ \begin{array}{c} 0 \\ 0 \\ \dots \\ 0 \\ a \end{array} \right] \left[ \begin{array}{c} 0 \\ 0 \\ \dots \\ b_1 \\ b_2 \end{array} \right] \dots \left[ \begin{array}{c} 0 \\ z_1 \\ \dots \\ z_{N-1} \end{array} \right] \\ \underbrace{\qquad\qquad\qquad}_{N-1} \end{array}$$

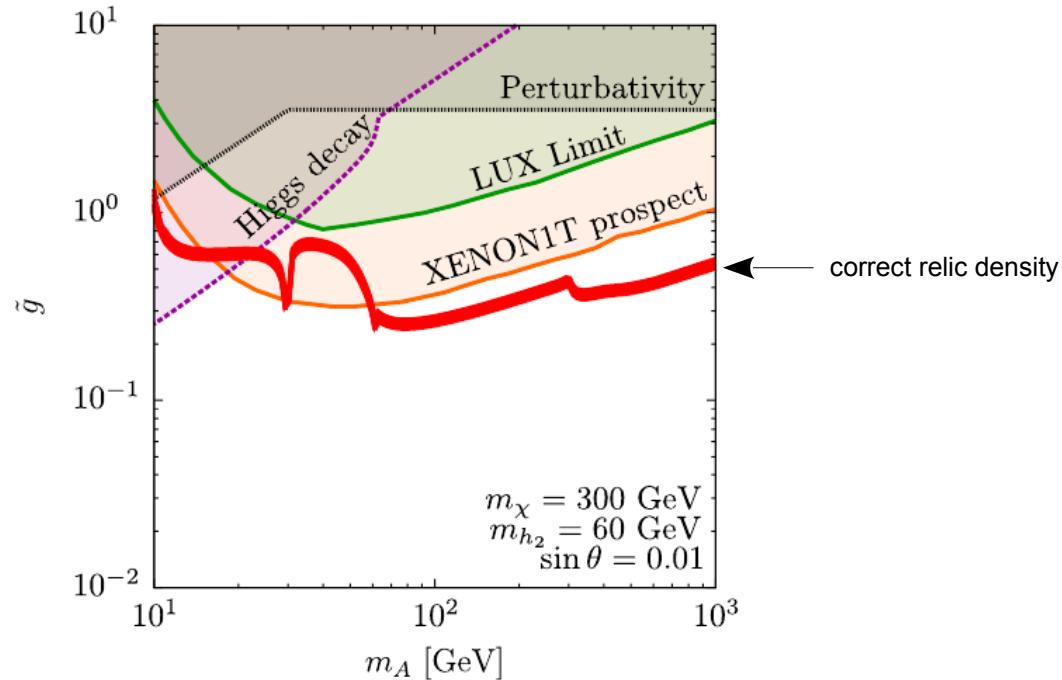
The scalar potential has an automatic global U(1) = overall phase multiplication

Unbroken U(1):

$$U = \underbrace{e^{-ix/N} \text{ diag } (e^{ix}, 1, \dots, 1)}_{SU(N)} \times \underbrace{e^{ix/N}}_{U(1)}$$

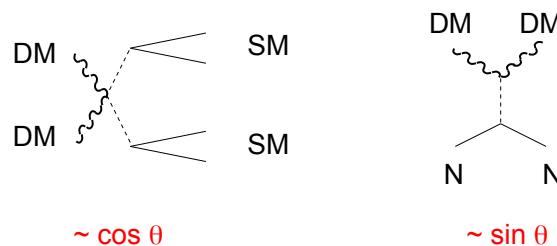
Gauge fields non-invariant under this U(1) cannot decay

## SU(3) example:



"Secluded" dark matter:

(à la Pospelov et al. '07)



## NB: non-WIMP options

- direct DM production by the inflaton (parametric resonance, decay, ...)
- freeze-in
- ...

initial state  $\rightarrow$  thermalization  $\rightarrow$  freeze-out

*Common approach:*

$$\left\{ \begin{array}{l} f(p) = \frac{1}{e^{\frac{E-\mu}{T}} - 1} \\ \mu \rightarrow 0 \end{array} \right. \quad \rightarrow \quad f(p) = e^{-(E-\mu)/T}$$

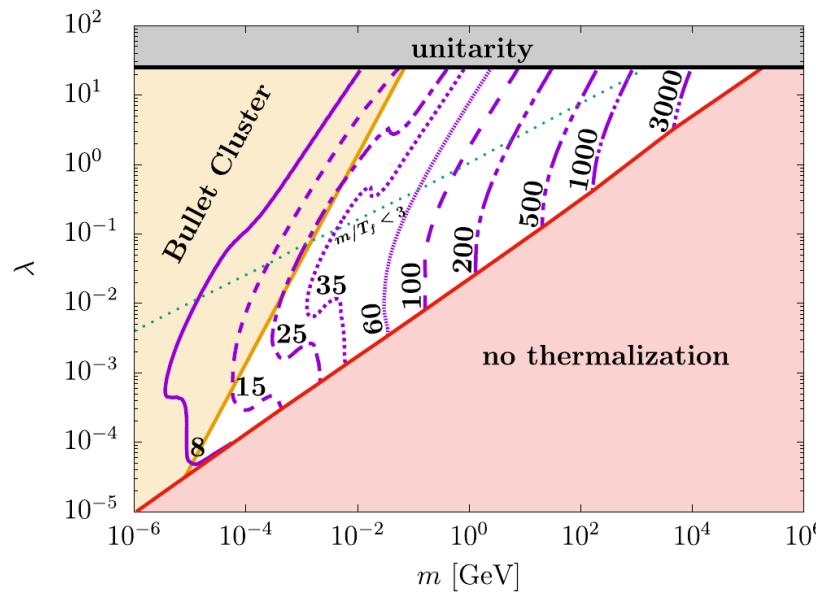
**Does not work in the relativistic regime  $T \gg m$  !**

Simplest model:

$$V = \frac{m^2}{2} S^2 + \frac{\lambda}{4!} S^4$$

**Relativistic analog of the Gelmini-Gondolo formula for the reaction rate:**

$$\Gamma_{2 \rightarrow 4} = \frac{4T}{\pi^4} \int_m^\infty dE E^3 \sqrt{E^2 - m^2} \int_0^\infty d\eta \frac{\sinh \eta}{e^{2(E \cosh \eta - \mu)/T} - 1} \ln \frac{\sinh \frac{E \cosh \eta + \sqrt{E^2 - m^2} \sinh \eta - \mu}{2T}}{\sinh \frac{E \cosh \eta - \sqrt{E^2 - m^2} \sinh \eta - \mu}{2T}} \sigma_{CM}(E, \eta)$$



includes BE final state factors

----- = boundary of relativistic freeze-out

8, 15, 25, ... =  $T_{SM} / T$

# The Higgs and inflation

*Expect on general grounds ( $\phi$  = inflaton):*

$$L = \lambda |H|^2 \phi^2 + \mu |H|^2 \phi$$

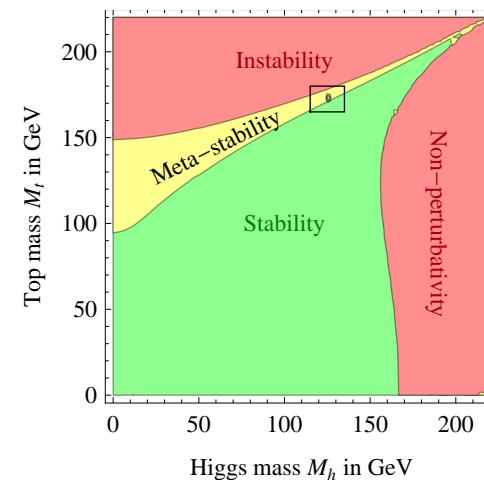
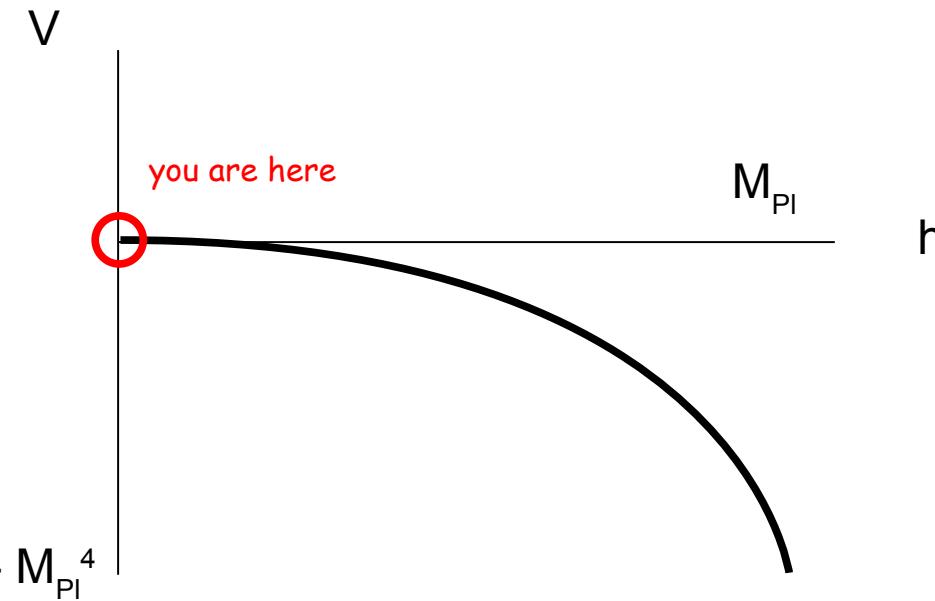


**both interaction and mixing!**



**Early Universe + colliders**

## (Apparent ) vacuum metastability:



Degassi et al.'12

OL, Westphal '12

- how did the Universe end up at  $h \sim 0$ ?
- why did it stay there during inflation ?

## "Minimal" solution:

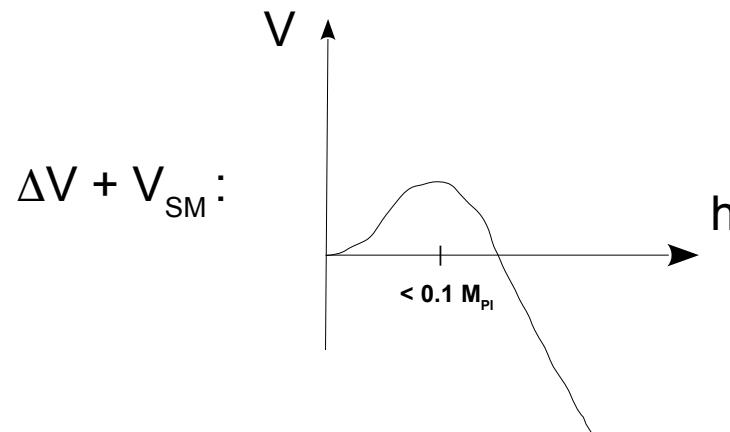
Higgs-inflaton coupling:

$$\Delta V = \frac{1}{2} \lambda_{h\phi} h^2 \phi^2$$

with  $\lambda_{h\phi} > 10^{-10}$

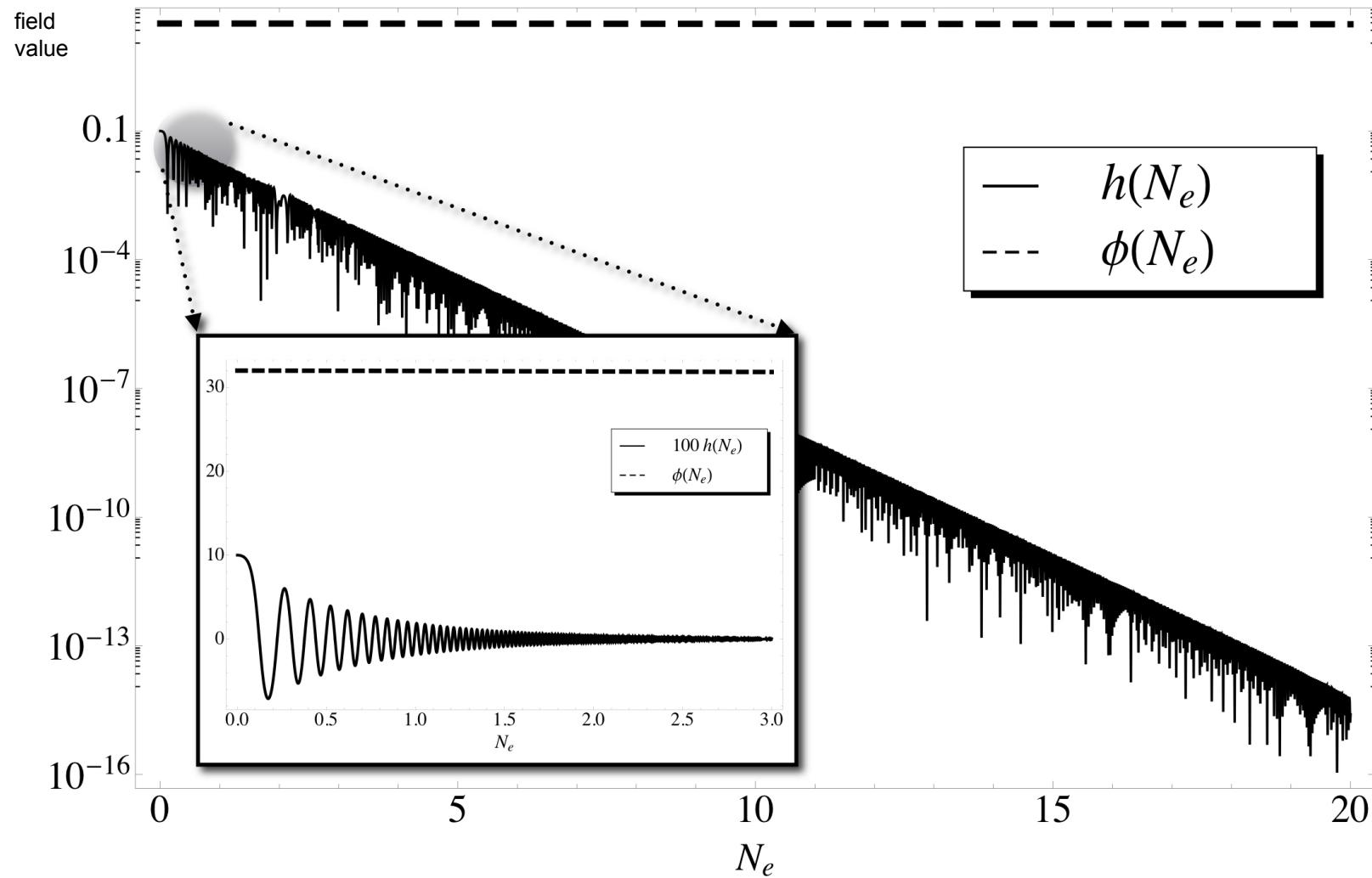
Espinosa, Giudice, Riotto '07  
Herranen, Markkanen, Nurmi, Rajantie '14

(or  $\xi R h^2$ )

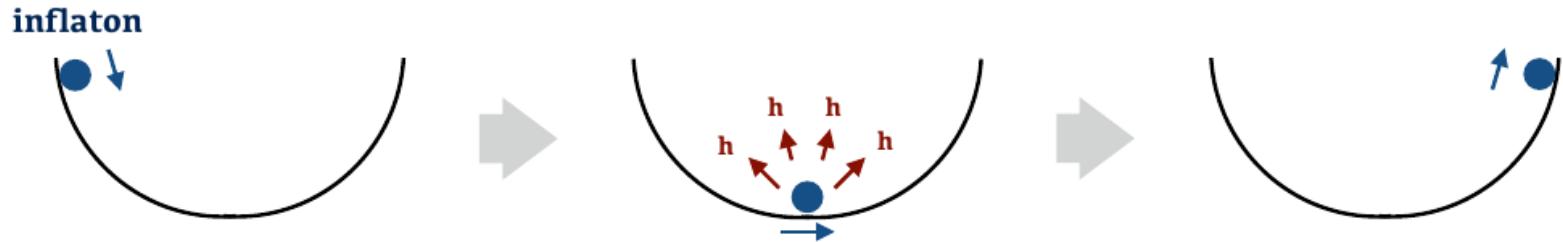


For all initial values of  $h$  up to  $0.1 M_{Pl}$ , the  $h$ -potential is convex  
(higher  $h$ -values  $\rightarrow$  Planckian density)

## Higgs/inflaton evolution (in $M_{\text{Pl}}$ ):



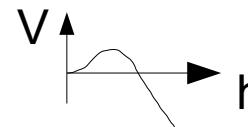
## Destabilization issue:



Kofman, Linde , Starobinsky '98

$$\Delta V = \frac{1}{2} \lambda_{h\phi} h^2 \phi^2 \quad \rightarrow \quad \text{parametric resonance}$$

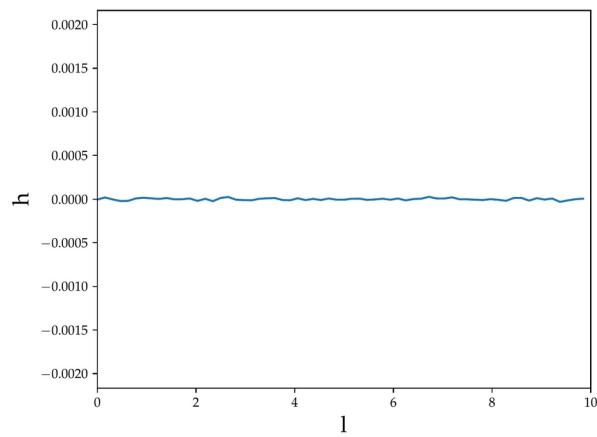
$$\langle h^2 \rangle \propto \text{Number of Higgs quanta}$$



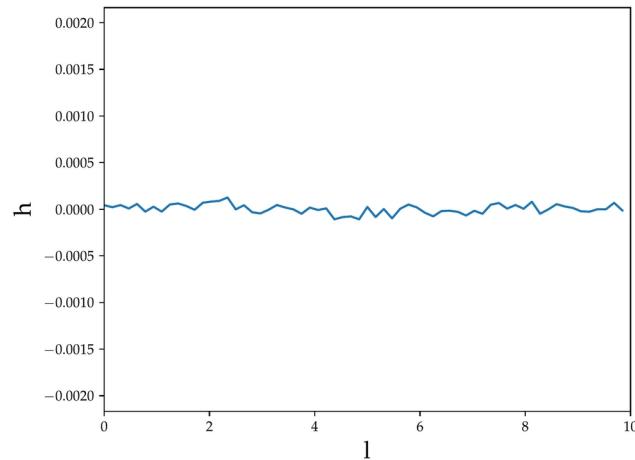
**destabilization for  $\lambda_{h\phi} > 10^{-8}$**

## Higgs field evolution (lattice results):

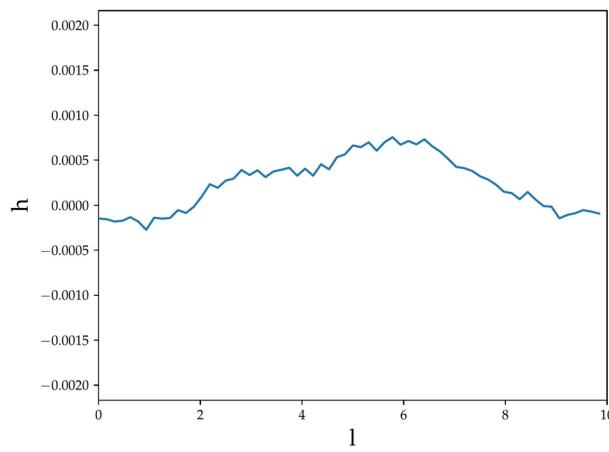
$t=0$



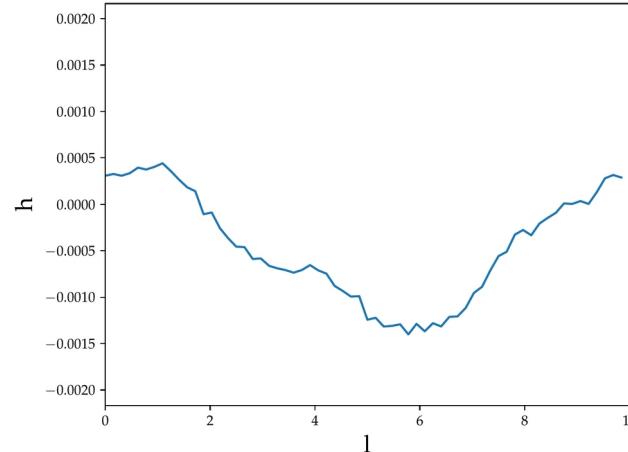
$t=15$



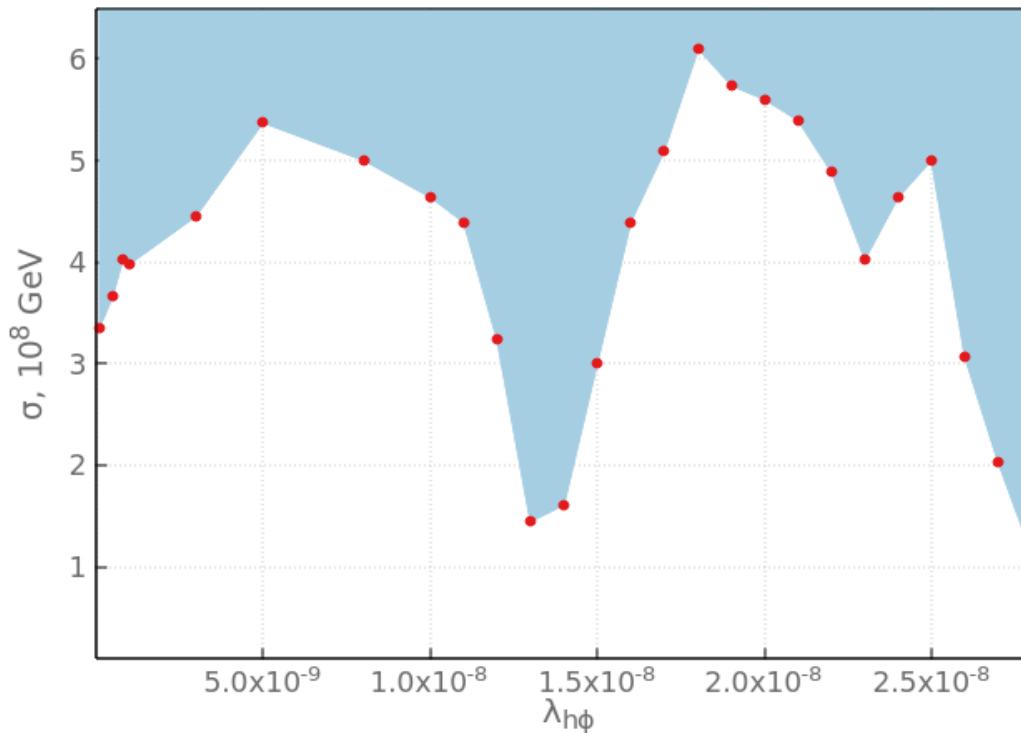
$t=27$



$t=30$



Bound on the trilinear coupling  $\phi h h$  :



## Stabilization via Higgs-inflaton mixing:

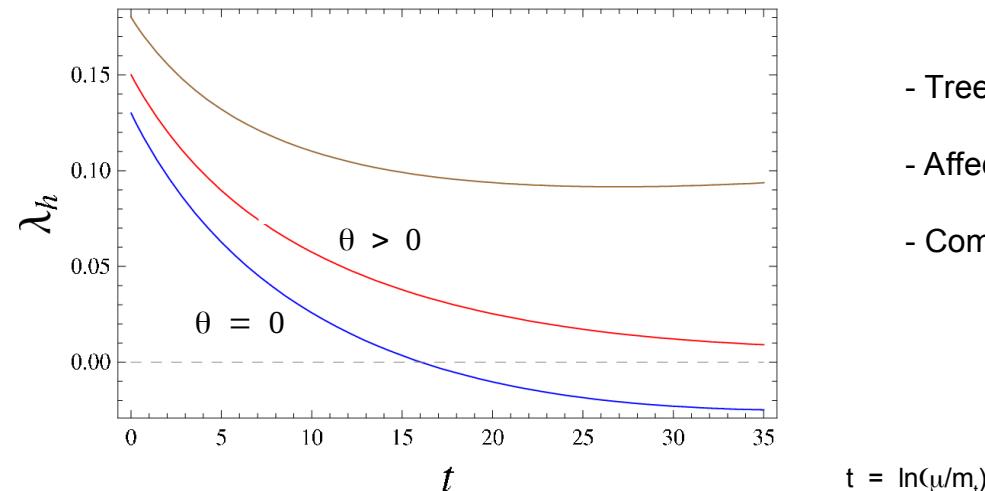
$$\Delta L \sim \phi^\dagger H^\dagger H$$

Two mass eigenstates  $h_{1,2}$  with mixing angle  $\theta$  :

$$2\lambda_h v^2 = m_1^2 \cos^2 \theta + m_2^2 \sin^2 \theta$$



$\lambda_h$  increases for  $m_2 > m_1$  !



- Tree level effect
- Affects boundary conditions
- Complete stability

## The Higgs+ inflaton lowest order terms :

$$S = \int d^4x \sqrt{-\hat{g}} \left[ \frac{1}{2} \Omega^2 \hat{R} - \frac{1}{2} \hat{g}^{\mu\nu} \partial_\mu \phi \partial_\nu \phi - \frac{1}{2} \hat{g}^{\mu\nu} \partial_\mu h \partial_\nu h - V(\phi, h) \right]$$

$$\Omega^2 = 1 + \xi_\phi \phi^2 + \xi_h h^2 ,$$

$$V(\phi, h) = \frac{\lambda_h}{4} h^4 - \frac{\mu_h^2}{2} h^2 + \frac{\lambda_{h\phi}}{2} h^2 \phi^2 + \sigma h^2 \phi + \frac{\lambda_\phi}{4} \phi^4 + \frac{b_3}{3} \phi^3 - \frac{\mu_\phi^2}{2} \phi^2 + b_1 \phi ,$$

Inflation is driven by  $\phi$  with  $\xi_\phi \gg 1$  ( à la Bezrukov-Shaposhnikov) :

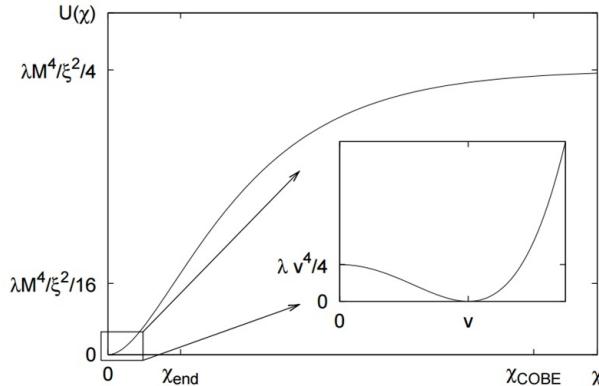


Fig. 1. Effective potential in the Einstein frame.

$$U(\chi) \simeq \frac{\lambda_\phi}{4\xi_\phi^2} \left( 1 - e^{-\sqrt{\frac{2}{3}}\chi} \right)^2$$

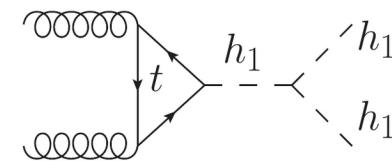
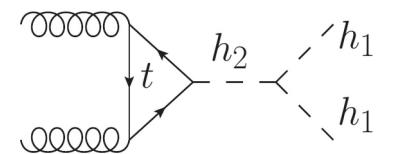
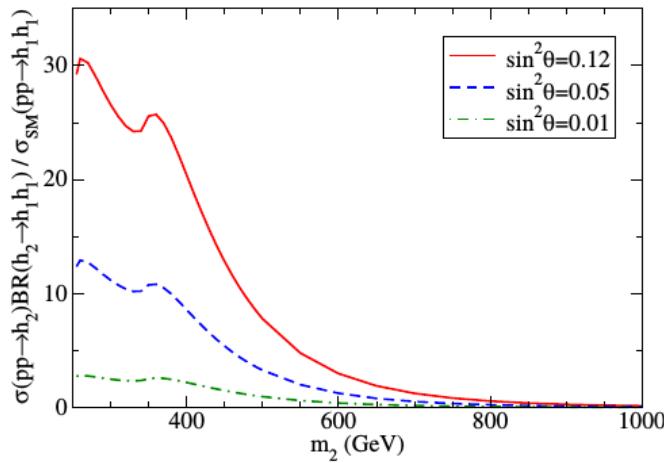
$$\chi \simeq \sqrt{\frac{3}{2}} \ln \xi_\phi \phi^2$$

- dimensionless couplings → inflation
- dimensionful parameters  $\ll M_{Pl}$  → low energy physics

## Inflaton search at LHC :

- universal Higgs coupling reduction
- heavy Higgs-like resonance
- resonant decay  $h_2 \rightarrow h_1 h_1$

Lewis, Sullivan '17



# Conclusion

- Higgs sector is special
- key to the hidden sector / DM / inflation
- to be tested at HL-LHC, DM detection, etc.