Extended Riemannian Geometry and Double Field Theory

Christian Sämann



School of Mathematical and Computer Sciences Heriot-Watt University, Edinburgh

Zagreb Workshop, 9.6.2017

Based on:

• arXiv:1611.02772 with Andreas Deser

Aim: Clarify mathematical structures behind Double Field Theory

Double Field Theory (DFT)

- target space formulations useful (SUSY YM, SUGRA)
- T-duality very interesting

Tseytlin, Siegel, Hull, Zwiebach, ...

• Success, e.g.: DFT action on \mathbb{R}^{2D} :

$$S = \int \mathrm{d}^{2D} x \, \mathrm{e}^{-2d} \, \left(\frac{1}{8} \mathcal{H}_{MN} \partial^M \mathcal{H}_{KL} \partial^N \mathcal{H}^{KL} - \frac{1}{2} \mathcal{H}_{MN} \partial^M \mathcal{H}_{KL} \partial^L \mathcal{H}^{KN} \right)$$

 $-2\partial^M d\partial^N \mathcal{H}_{MN} + 4\mathcal{H}_{MN}\partial^M d\partial^N d$

 \mathcal{H}_{MN} : generalized metric, d: DFT dilaton

- Reproduces SUGRAs previously not derived from string theory
- sensible? strange truncation of string modes? seems to work

DFT

- Massless string modes: metric g, 2-form B, dilaton ϕ .
- g and B transform differently, combine in generalized metric:

$$\mathcal{H}_{MN} = \begin{pmatrix} g_{\mu\nu} - B_{\mu\kappa}g^{\kappa\lambda}B_{\lambda\nu} & B_{\mu\kappa}g^{\kappa\nu} \\ -g^{\mu\kappa}B_{\kappa\nu} & g^{\mu\nu} \end{pmatrix} \quad \text{on} \quad T^*M \oplus TM$$

- Manifest T-duality: double space, coords.: $x^M = (x^\mu, x_\mu)$
- Level matching: Fields ϕ in DFT satisfy $\Box \phi = 0$
- Algebra of fields: strong section condition:

$$\partial^M \phi \partial_M \psi = 0$$

• Generalized Lie derivative wrt. $X = X^{\mu} \frac{\partial}{\partial x^{\mu}} + X_{\mu} \frac{\partial}{\partial x_{\mu}}$:

 $\hat{\mathcal{L}}_X \mathcal{H}_{MN} = X^P \partial_P \mathcal{H}_{MN} + (\partial_M X^P - \partial^P X_M) \mathcal{H}_{PN} + (\partial_N X^P - \partial^P X_N) \mathcal{H}_{MP}$

• $\hat{\mathcal{L}}_X$ form Lie algebra, but not representation of Lie algebra $\hat{\mathcal{L}}_X \hat{\mathcal{L}}_Y - \hat{\mathcal{L}}_Y \hat{\mathcal{L}}_X = \hat{\mathcal{L}}_{\mu_2(X,Y)}$,

• $\mu_2(X,Y)$ known as C-bracket (actual interpretation later)

Issues to address in a mathematical formulation:

- Issue 1: Algebraic structure behind C-bracket
- Issue 2: Section condition is rather invidious:

 $\partial^M \phi \ \partial_M \psi = 0$

- Issue 3: Section condition is often too strong
- Issue 4: Reasonable version of Doubled Riemannian Geometry

$$\widehat{\Gamma}_{MNK} = -2(P\partial_M P)_{[NK]} - 2(\bar{P}_{[N}{}^P \bar{P}_{K]}{}^Q - P_{[N}{}^P P_{K]}{}^Q)\partial_P P_{QM} + \frac{4}{D-1}(P_{M[N}P_{K]}{}^Q + \bar{P}_{M[N}\bar{P}_{K]}{}^Q)(\partial_Q d + (P\partial^P P)_{[PQ]})$$

- Issue 5: Global formulation of DFT
- Issue 6: Same for Exceptional field theory (M-theory analogue)

- Massless string modes: metric g, 2-form B, dilaton ϕ .
- Recall: *B*-field belongs to gerbe/categorified principal bundle
- Generalized Geometry: Courant/Lie 2-algebroid $TM\oplus T^*M$
- Captures Lie 2-algebra of symmetries (gauge+diffeos)
- In particular: Courant/Dorfman brackets
- Picture allows for an extension to pre-Lie 2-algebroids
- Lie 2-algebra with C- and D-brackets
- Fully algebraic and coordinate independent picture
- Know how to twist pre-Lie 2-algebroids
- Potential global picture (work in progress)

Part I : Generalized Geometry

- Gerbes lead to Courant algebroids
- Higher Lie algebroids/algebras as NQ-manifolds
- Symmetries from higher symplectic Lie algebroids

Ingredients: massless string modes: metric g, 2-form B, dilaton ϕ .

• Well-known: *B* belongs to connective structure of a gerbe Gawedzki 1987, Freed&Witten 1999

Gerbes - Intuitive Picture

- Take your favorite definition of principal bundles
- Smooth Manifold \longrightarrow smooth categories $\mathcal{M}_1 \rightrightarrows \mathcal{M}_0$
- Smooth morphisms between manifolds \longrightarrow smooth functors
- Structure group \longrightarrow special category with products
- There are obvious categorified notions of: fibration, surjection, diffeomorphism, free, transitive action
- Result: principal 2-bundles
- Gerbe: principal 2-bundle for $U(1) \rightrightarrows *$

Recall: Čech Description of a Gerbe Generalized Geometry provides a nice geometric description of the *B*-field.

Cover/Surjective submersion

Manifold M, cover $Y = \sqcup_i U_i \twoheadrightarrow M$, $Y^{[2]} = \sqcup_{i,j} U_i \cap U_j$, ...

Principal U(1)-bundle over M

 $U(1) \to P \to M$, $g_{ij}g_{jk} = g_{ik}$, $A_i = g_{ij}^{-1}(A_j + d)g_{ij}$

Symmetries: U(1) gauge symmetry and diffeomorphisms.

Ordinary geometry: "Local" Atiyah Algebroid for U(1)-bundle $0 \rightarrow M \times i\mathbb{R} \xrightarrow{\hookrightarrow} TM \oplus i\mathbb{R} \xrightarrow{\operatorname{pr}} TM \rightarrow 0$ • Gauge potential $A: TM \rightarrow i\mathbb{R}$ is defined by section of pr • Infinitesimal symm. (diffeos+gauge): sections of $TM \oplus i\mathbb{R}$

Generalized Geometry: Courant algebroid for U(1)-gerbe

$$0 \to T^*M \xrightarrow{\hookrightarrow} TM \oplus T^*M \xrightarrow{\operatorname{pr}} TM \to 0$$

- $g + B : TM \to T^*M$ again from section of pr
- ullet Infinitesimal symm. (diffeos+gauge): sections of $TM\oplus T^*M$



N-manifolds, NQ-manifold

 $\bullet~\mathbb{N}_0\text{-}\mathsf{graded}$ manifold with coordinates of degree $0,1,2,\ldots$



- NQ-manifold: vector field Q of degree 1, $Q^2 = 0$
- Physicists: think ghost numbers, BRST charge, SFT

Examples:

- Tangent algebroid T[1]M, $\mathcal{C}^{\infty}(T[1]M) \cong \Omega^{\bullet}(M)$, Q = d
- Lie algebra $\mathfrak{g}[1]$, coordinates ξ^a of degree 1:

$$Q = -\frac{1}{2} f^c_{ab} \xi^a \xi^b \frac{\partial}{\partial \xi^c} \quad , \quad \text{Jacobi identity} \Leftrightarrow \ Q^2 = 0$$

Lie *n*-algebroid: $M_0 \leftarrow M_1 \leftarrow M_2 \leftarrow \ldots \leftarrow M_n \leftarrow * \leftarrow * \leftarrow \ldots$ Lie *n*-algebra or *n*-term L_∞ -algebra:

 $* \leftarrow M_1 \leftarrow M_2 \leftarrow \ldots \leftarrow M_n \leftarrow * \leftarrow * \leftarrow \ldots$

Important example: Lie 2-algebra

- Graded vector space: $W[1] \leftarrow V[2]$
- Coordinates: w^a of degree 1 on W[1], v^i of degree 2 on V[2]• Most general vector field Q of degree 1:

$$Q = -m_{i}^{a}v^{i}\frac{\partial}{\partial w^{a}} - \frac{1}{2}m_{ab}^{c}w^{a}w^{b}\frac{\partial}{\partial w^{c}} - m_{ai}^{j}w^{a}v^{i}\frac{\partial}{\partial v^{j}} - \frac{1}{3!}m_{abc}^{i}w^{a}w^{b}w^{c}\frac{\partial}{\partial v^{i}}$$

• Induces "brackets"/"higher products":

 $\mu_1(\tau_i) = m_i^a \tau_a , \quad \mu_2(\tau_a, \tau_b) = m_{ab}^c \tau_c , \quad \dots , \quad \mu_3(\tau_a, \tau_b, \tau_c) = m_{abc}^i \tau_i$

- $Q^2 = 0 \Leftrightarrow$ Homotopy Jacobi identities, e.g. $\mu_1(\mu_1(-)) = 0$
- Failure of Jacobi identity: $\mu_2(x, \mu_2(y, z)) + \ldots = \mu_1(\mu_3(x, y, z))$



• \mathbb{N}_0 -graded manifold, vector field Q of degree 1, $Q^2 = 0$



• symplectic NQ-manifold: ω nondegenerate, closed, $\mathcal{L}_Q \omega = 0$

Examples of symplectic NQ-manifolds:

Metric Lie algebra:

 $(\mathfrak{g}[1], \omega = g_{\alpha\beta} \mathrm{d}\xi^{\alpha} \wedge \mathrm{d}\xi^{\beta}), \quad Q = -\frac{1}{2} f_{ab}^{c} \xi^{a} \xi^{b} \frac{\partial}{\partial \xi^{c}}$

• Symplectic manifolds (symplectic Lie 0-algebroids):

 (M,ω) , Q = 0• Poisson manifolds (symplectic Lie 1-algebroids):

 $T^*[1]M$, $\omega = \mathrm{d}p_\mu \wedge \mathrm{d}x^\mu$, $Q = \pi^{\mu\nu} p_\mu \frac{\partial}{\partial x^\nu}$

• Courant algebroids (symplectic Lie 2-algebroids, next slide...)

The Courant Algebroid as a Symplectic NQ-Manifold The key geometric structure behind generalized geometry is locally $T^*[2]T[1]M$.

The symplectic NQ-manifold $T^*[2]T[1]M$

Local description, choose $M = \mathbb{R}^D$ as base manifold.

 $T^{*}[2]T[1]M = \mathbb{R}^{D} \oplus \mathbb{R}^{D}[1] \oplus \mathbb{R}^{D}[1] \oplus \mathbb{R}^{D}[2]$ $x^{\mu} \quad \xi^{\mu} \quad \zeta_{\mu} \quad p_{\mu}$ $Q = \xi^{\mu}\frac{\partial}{\partial x^{\mu}} + p_{\mu}\frac{\partial}{\partial \zeta_{\nu}} \quad \omega = \mathrm{d}x^{\mu} \wedge \mathrm{d}p_{\mu} + \mathrm{d}\xi^{\mu} \wedge \mathrm{d}\zeta_{\mu}$

• functions on
$$T^*[2]T[1]M$$
 of degree 0: functions on M

- functions on $T^*[2]T[1]M$ of degree 1: sections of $TM \oplus T^*M$
- Poisson bracket induced by ω , pairing on degree 1 functions:

 $\{X + \alpha, Y + \beta\} = \iota_X \alpha + \iota_Y \beta$

The L_{∞} -Algebra of a Symplectic L_{∞} -algebroid Every symplectic L_{∞} -algebroid comes with an L_{∞} -algebra.

An underappreciated/widely unknown fact: All symplectic Lie *n*-algebroids \mathcal{M} come with Lie *n*-algebra $L(\mathcal{M})$: $\mathcal{C}_{0}^{\infty}(\mathcal{M}) \xrightarrow{Q} \mathcal{C}_{1}^{\infty}(\mathcal{M}) \xrightarrow{Q} \dots \xrightarrow{Q} \mathcal{C}_{n-2}^{\infty}(\mathcal{M}) \xrightarrow{Q} \mathcal{C}_{n-1}^{\infty}(\mathcal{M})$ $\mathsf{L}_{n-1}(\mathcal{M}) \xrightarrow{Q} \mathsf{L}_{n-2}(\mathcal{M}) \xrightarrow{Q} \dots \xrightarrow{Q} \mathsf{L}_{1}(\mathcal{M}) \xrightarrow{Q} \mathsf{L}_{0}(\mathcal{M})$ $\mu_1(\ell) = \begin{cases} 0 & \ell \in \mathcal{C}_{n-1}^{\infty}(\mathcal{M}) = \mathsf{L}_0(\mathcal{M}) \\ Q\ell & \text{else} \end{cases}$ $\mu_2(\ell_1,\ell_2) = \frac{1}{2} \left(\{ \delta\ell_1,\ell_2 \} \pm \{ \delta\ell_2,\ell_1 \} \right) , \quad \delta(\ell) = \begin{cases} Q\ell & \ell \in \mathsf{L}_0(\mathcal{M}) \\ 0 & \text{else} \end{cases}$ $\mu_3(\ell_1, \ell_2, \ell_3) = -\frac{1}{12} (\{\{\delta\ell_1, \ell_2\}, \ell_3\} \pm \dots)$

Roytenberg, Rogers, Fiorenza/Manetti, Getzler

- "Derived brackets" (Kosmann-Schwarzbach, Voronov, ...)
- Important class of examples next
- Cartan calculus (⇒ Andreas' talk)

Associated Algebraic Structures of Courant Algebraid The key geometric structure behind generalized geometry is locally $T^*[2]T[1]M$.

• $\mathcal{M} = T^*[2]T[1]\mathbb{R}^D = \mathbb{R}^D \oplus \mathbb{R}^{2D}[1] \oplus \mathbb{R}^D[2], Q, \omega, \{-,-\}$

Graded vector space:

 $\begin{array}{ll} \mathcal{C}_0^{\infty}(\mathcal{M}) \to \mathcal{C}_1^{\infty}(\mathcal{M}) & = & \mathcal{C}^{\infty}(\mathbb{R}^D) \to \Gamma(T\mathbb{R}^D \oplus T^*\mathbb{R}^D) \\ & = & \mathsf{L}_1(\mathcal{M}) \to \mathsf{L}_0(\mathcal{M}) \end{array}$

• Derived brackets from differential $\delta(\ell) = \begin{cases} Q\ell & \ell \in L_0(\mathcal{M}) \\ 0 & \text{else} \end{cases}$

- Leibniz algebra from Dorfman bracket $\nu_2(\ell_1,\ell_2) = \{\delta\ell_1,\ell_2\}$
- Lie 2-algebra from Courant bracket $\mu_2(\ell_1, \ell_2) = \{\delta \ell_{[1}, \ell_{2]}\}$
- This is the symmetry Lie 2-algebra of corresponding gerbe!



Part of Larger Family: GR + n-Form Potentials 16/2 The above picture is part of a larger story, involving GR coupled to *n*-form potentials.

More generally: Couple GR to n-form gauge potential

Vinogradov algebroids $\mathcal{V}_n(M)$

- Locally as a vector bundle: $\mathcal{V}_n(M) := T^*[n]T[1]M$
- coords: $(x^{\mu},\xi^{\mu},\zeta_{\mu},p_{\mu})$ of degrees (0,1,n-1,n)
- Homological vector field: $Q = \xi^{\mu} \frac{\partial}{\partial x^{\mu}} + p_{\mu} \frac{\partial}{\partial \zeta_{\mu}}$
- Symplectic form: $\omega = dx^{\mu} \wedge dp_{\mu} + d\xi^{\mu} \wedge d\zeta_{\mu}$
- Geometric picture: principal n-bundle or n 1-gerbe
- diffeos+gauge: degree n-1 functions of $\mathcal{V}_n(M)$.
- Higher Courant/Dorfman brackets
- Full Symmetrie Lie *n*-algebra: Lie *n*-algebra of $\mathcal{V}_n(M)$

How to extend this to Double Field Theory?

Part II : Double Field Theory / Extended Geometry

- Generalized NQ-manifolds
- Construct example suitable for DFT
- Weakened section condition as algebraic relation
- Comments on global picture

Weakening the Definition of Symplectic NQ-Manifolds 18/26 There are two properties that offer themselves to a weakening.

Requirements

- Geometry built from (doubled) spacetime
- Symmetry Lie *n*-algebra structure from derived brackets
- Reduction to Vinogradov algebroids of Generalized Geometry
- a few further points, related to global picture

Need to keep: graded vector bundle, symplectic form (?), |Q| = 1

Note: $Q^2 = 0$ not necessary everywhere for L_{∞} -algebra:

$$\mathsf{L}_{n-1}(\mathcal{M}) \xrightarrow{Q} \mathsf{L}_{n-2}(\mathcal{M}) \xrightarrow{Q} \dots \xrightarrow{Q} \mathsf{L}_{1}(\mathcal{M}) \xrightarrow{Q} \mathsf{L}_{0}(\mathcal{M}) \xrightarrow{0} 0$$

Instead: something like $\{Q^2-,-\}=0$.

Definition: Symplectic pre-NQ-Manifold of degree n

Symplectic N-manifold (\mathcal{M}, ω) of degree n, i.e. $|\omega| = n$, compatible vector field Q of degree 1, i.e. |Q| = 1 and $\mathcal{L}_Q \omega = 0$.

Definition: L_{∞} -structure

A subset $L(\mathcal{M})$ of the functions $\mathcal{C}^{\infty}(\mathcal{M})$ such that the derived brackets close and form an L_{∞} -algebra.

Theorem

$$\begin{split} (\mathcal{M}, \omega) \text{ of degree 2.} \\ \mathsf{L}(\mathcal{M}) &= \mathsf{L}_1(\mathcal{M}) \oplus \mathsf{L}_0(\mathcal{M}) \subset \mathcal{C}^\infty(\mathcal{M}) \text{, derived brackets close.} \\ \mathsf{L}(\mathcal{M}) \text{ is } L_\infty \text{-structure iff for all } f, g \in \mathsf{L}_1(\mathcal{M}), X, Y, Z \in \mathsf{L}_0(\mathcal{M}) \text{:} \\ & \{Q^2 f, g\} + \{Q^2 g, f\} = 0 , \quad \{Q^2 X, f\} + \{Q^2 f, X\} = 0 \\ & \{\{Q^2 X, Y\}, Z\}_{[X,Y,Z]} = 0 \end{split}$$

Application to DFT: Built pre-NQ-manifold The pre-NQ-manifold for DFT from a kind of polarization.

- Start: $M=\mathbb{R}^D$, double: T^*M , coords. $x^M=(x^\mu,x_\mu)$
- $\mathcal{V}_2(T^*M) = T^*[2]T[1](T^*M)$, coords. $(x^M, \xi^M, \zeta_M, p_M)$
- $\boldsymbol{\omega} = \mathrm{d}x^M \wedge \mathrm{d}p_M + \mathrm{d}\xi^M \wedge \mathrm{d}\zeta_M$, $\boldsymbol{Q} = \sqrt{2} \left(\xi^M \frac{\partial}{\partial x^M} + p_M \frac{\partial}{\partial \zeta_M} \right)$
- Crucial to GenGeo/DFT: Local symmetry group O(D, D)
- $\bullet~\mbox{Reduce}$ structure group ${\rm GL}(2D,\mathbb{R})$ to ${\rm O}(D,D)$ by introducing

$$\eta_{MN} = \eta^{MN} = \left(\begin{array}{cc} 0 & \mathbb{1} \\ \mathbb{1} & 0 \end{array}\right)$$

New coordinates:

$$\boldsymbol{\theta}^{M} = \frac{1}{\sqrt{2}} (\boldsymbol{\xi}^{M} + \boldsymbol{\eta}^{MN} \boldsymbol{\zeta}_{N}) \qquad \boldsymbol{\beta}^{M} = \frac{1}{\sqrt{2}} (\boldsymbol{\xi}^{M} - \boldsymbol{\eta}^{MN} \boldsymbol{\zeta}_{N})$$

• Polarize by putting $\beta^M \stackrel{!}{=} 0$:

$$Q = \theta^M \frac{\partial}{\partial x^M} + p_M \eta^{MN} \frac{\partial}{\partial \theta^N} , \quad Q^2 = p_M \eta^{MN} \frac{\partial}{\partial x^N} \neq 0$$

• We obtain a symplectic pre-NQ-manifold, $\mathcal{E}_2(M)$.

Application to DFT: Weakened Section Condition This framework is readily applied to reproduce DFT's strong section condition.

$$\begin{aligned} \mathcal{E}_{2}(\mathbb{R}^{D}) &= \mathbb{R}^{2D} \oplus \mathbb{R}^{2D}[1] \oplus \mathbb{R}^{2D}[2] \\ & x^{M} & \theta^{M} & p_{M} \end{aligned}$$
$$Q &= \theta^{M} \frac{\partial}{\partial x^{M}} + p_{M} \eta^{MN} \frac{\partial}{\partial \theta^{N}} \quad \boldsymbol{\omega} = \mathrm{d} x^{M} \wedge \mathrm{d} p_{M} + \frac{1}{2} \eta_{MN} \mathrm{d} \theta^{M} \wedge \mathrm{d} \theta^{N} \end{aligned}$$

Proposition

Let $L = L_0 \oplus L_1$ be L_{∞} -structure on $\mathcal{E}_2(\mathbb{R}^D)$, $f, g \in L_1$ and $X = X_M \theta^M, Y = Y_M \theta^M, Z = Z_M \theta^M \in L_0$. Then:

$$\{Q^{2}f,g\} + \{Q^{2}g,f\} = 2\partial^{M}f \ \partial_{M}g = 0$$

$$\{Q^{2}X,f\} + \{Q^{2}f,X\} = 2\partial^{M}X \ \partial_{M}f = 0$$

$$\{\{Q^{2}X,Y\},Z\}_{[X,Y,Z]} = 2\theta^{L} ((\partial^{M}X_{L})(\partial_{M}Y^{K})Z_{K})_{[X,Y,Z]} = 0$$

Note: Fulfilled for $\partial^M \phi \ \partial_M \psi = 0 \Rightarrow$ Weakened section condition

The picture we have:

- Symmetries of DFT are described by L_∞ -structure on $\mathcal{E}_2(\mathbb{R}^D)$
- Extended Geometry $\mathcal{E}_2(\mathbb{R}^D)$: polarization of GenGeo
- Reduce Ext. Geo. \rightarrow GenGeo: Choice of L_{∞} -structure
- Example:

$$\mathsf{L} = \left\{ F \in \mathcal{C}^{\infty}(\mathcal{E}_{2}(\mathbb{R}^{D})) | \frac{\partial}{\partial x_{\mu}} F = 0 \right\}$$
$$Q = \theta^{\mu} \frac{\partial}{\partial x^{\mu}} + p_{\mu} \frac{\partial}{\partial \theta^{\mu}} \qquad \omega = \mathrm{d}x^{\mu} \wedge \mathrm{d}p_{\mu} + \mathrm{d}\theta^{\mu} \wedge \mathrm{d}\theta_{\mu}$$

 \Rightarrow Vinogradov algebroid $\mathcal{V}_2(\mathbb{R}^D)$ of Generalized Geometry

Observations:

- Action of Lie algebra \mathfrak{g} on manifold M: hom. $\mathfrak{g} \to \mathfrak{X}(M)$
- N-manifold \mathcal{M} , $\mathfrak{X}(\mathcal{M})$ is \mathbb{N}_0 -graded Lie algebra, L_∞ -algebra

Definition

Action of L_{∞} -algebra L on manifold \mathcal{M} : L_{∞} -morph. L $\rightarrow \mathfrak{X}(\mathcal{M})$.

Note: $\mathcal{M} = \mathcal{V}_2(M)$ only encodes forms, not symmetric tensors.

Extension of Poisson bracket

$$\{-,-\}: \ \mathcal{C}^{\infty}(\mathcal{M}) \times T(\mathcal{M}) \to T(\mathcal{M})$$

 $\{f, g \otimes h\} := \{f, g\} \otimes h + (-1)^{(n-|f|)|g|}g \otimes \{f, h\}$

Definition: Extended tensors

Let L be L_{∞} -structure on \mathcal{M} . Extended tensors are elts. of $T(\mathcal{M})$ such that elements of L act on it via $X \triangleright t := \{\delta X, t\}$.

Comments on Global Picture: Generalized Geometry The global picture seems to be in reach now.

- Surjective submersion $Y \twoheadrightarrow M$, $Y^{[2]} := Y \times_M Y$, $Y^{[n]} := ...$
- For example, $Y = \sqcup_i U_i$, $Y^{[2]} = \sqcup_{i,j} U_i \cap U_j$
- ${\, \bullet \, }$ Gerbe is a principal U(1)-bundle over ${\it Y}^{[2]}$ + data over ${\it Y}^{[3]}$
- Trivial gerbe: Y = M, $Y^{[2]} = M$, $\mathscr{G} = M \times U(1) = M \times S^1$.
- T-duality on $M \times S^1$: trivial gerbe, H = dB globally.

Non-trivial gerbe in Generalized Geometry

- $\mathcal{V}_2(M) = T^*[2]T[1]M$ and $Q = \xi^{\mu} \frac{\partial}{\partial x^{\mu}} + p_{\mu} \frac{\partial}{\partial \zeta_{\mu}}$ only locally
- Assume nontrivial gerbe $H, B_{(i)}, A_{(ij)}, h_{(ijk)}$
- Hitchin's Generalized tangent bundle over patches i, j: $(X_{(i)}, \Lambda_{(i)}) = (X_{(j)}, \Lambda_{(j)} + \iota_X dA_{(ij)})$
- $Q = \xi^{\mu} \frac{\partial}{\partial x^{\mu}} + p_{\mu} \frac{\partial}{\partial \zeta_{\mu}} \frac{1}{3!} (\frac{\partial}{\partial x^{\nu}} H) \frac{\partial}{\partial p_{\nu}} + \frac{1}{2!} H_{\nu \mu_{1} \mu_{2}} \xi^{\mu_{1}} \xi^{\mu_{2}} \frac{\partial}{\partial \zeta_{\nu}}$

- ${\ \bullet \ } L_\infty{\ } {\ }$ algebra of symmetries of DFT acts as Lie algebra
- Lie algebra can be integrated Hohm, Zwiebach, 2012
- Proposal: Patch local descriptions by finite DFT symmetries Berman, Cederwall, Perry, 2014
- Papadopoulos 2014: This only works for trivial gerbes
- No surprise, DFT reduces to GenGeo, where we need to twist!
- Need to twist C-/D-bracket, just as in GenGeo.
- Twist can be defined and studied in our framework.
- We recover twists of Generalized Geometry as special cases.
- Integrate twisted action
- Global picture: Patch together with twisted transformations!
- Understand all this \Rightarrow potentially win Berman medal

Summary:

- ✓ Full algebraic and geometric picture for local DFT
- ✓ Picture is extension from GR coupled to n-form fields
- \checkmark weakened section condition from algebra
- ✓ twist of symmetry Lie 2-algebra
- $\checkmark\,$ Initial studies of global picture and Riemannian Geometry More:
 - \Rightarrow Andreas' talk this afternoon

Soon to come:

- ▷ Heteroric DFT
- Exceptional Field Theory (M-theory)
- Global Picture

Extended Riemannian Geometry and Double Field Theory

Christian Sämann



School of Mathematical and Computer Sciences Heriot-Watt University, Edinburgh

Zagreb Workshop, 9.6.2017